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Degertekin et al.

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(54) **METHOD AND APPARATUS FOR THE
ULTRASONIC ACTUATION OF THE
CANTILEVER OF A PROBE-BASED
INSTRUMENT**

5,723,775 A 3/1998 Watanabe et al.
5,742,377 A 4/1998 Minne et al.
5,908,981 A 6/1999 Atalar et al.
6,389,885 B1 * 5/2002 Arnold et al. 73/105

OTHER PUBLICATIONS

Han et al., *A Magnetically Driven Oscillating Probe Microscope for Operation in Liquids*, American Institute of Physics [S0003-6951(96)05152-2], pp. 4111-4113 (1996).

Putman et al., *Tapping Mode Atomic Force Microscopy in Liquid*, Appl. Phys. Lett. 64 (18), pp. 2454-2456 (May 2, 1994).

Lawrence A. Bottomley, *Scanning Probe Microscopy*, Anal. Chem. 1998, vol. 70, No. 12, Jun. 15, 1998, pp. 425R-475R.

Ratneshwar Lal and Scott A. John, *Biological applications of atomic force microscopy*, The American Physiological Society, 1994, pp. C1-C21.

(List continued on next page.)

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(52) **U.S. Cl.** **73/661; 73/652; 73/105**

(58) **Field of Search** 73/652, 649, 579,
73/105, 661; 438/52

(56) **References Cited**

U.S. PATENT DOCUMENTS

4,977,779 A 12/1990 Karaki et al.
5,266,801 A 11/1993 Elings et al.
RE34,489 E 12/1993 Hansma et al.
5,319,977 A 6/1994 Quate et al.
5,412,980 A 5/1995 Elings et al.
5,436,448 A 7/1995 Hosaka et al.
5,517,280 A 5/1996 Quate
5,670,712 A 9/1997 Cleveland et al.

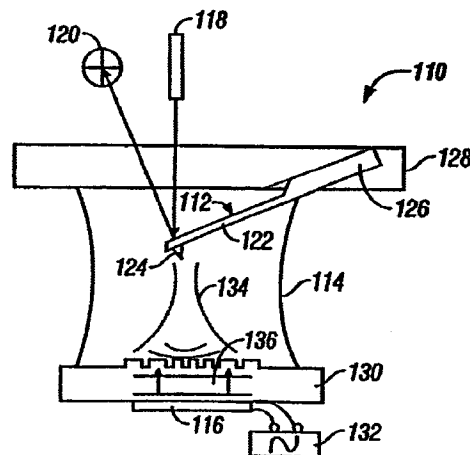
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Newholm Stein & Gratz S.C.

(57) **ABSTRACT**

The cantilever of a probe-based metrology instrument such as an AFM is deflected by directing a beam of ultrasonic energy at the cantilever to apply ultrasonically generated acoustic radiation pressure to the cantilever. The energy is generated by an ultrasonic actuator such as a ZnO transducer driven by a power source such as an RF signal generator. The transmitted beam preferably is shaped by focusing, collimation, or the like so that it impinges at least primarily on a region of interest of the cantilever such as the free end. The ultrasonic actuator produces a much better controlled force on the cantilever than can be achieved through the use of a traditional piezoelectric actuator and, accordingly, produces a response free of spurious effects (at least when the cantilever is operating in liquid). It also has a frequency bandwidth in the MHz range.

24 Claims, 18 Drawing Sheets



OTHER PUBLICATIONS

- Marco Tortonese and Michael Kirk, *Characterization of application specific probes for SPMs*, SPIE vol. 3009 (undated), pp. 53–60.
- J.P. Cleveland, S. Manne, D. Bocek, and P.K. Hansma, *A nondestructive method for determining the spring constant of cantilevers for scanning force microscopy*, Rev. Sci. Instrum. 64(2), Feb. 1993, pp. 403–405.
- John E. Sader, James W.M. Chon, and Paul Mulvaney, *Calibration of rectangular atomic force microscope cantilevers*, American Institute of Physics, (1999), pp. 3967–3969.
- K.B. Crozier, G.G. Yaralioglu, F.L. Degertekin, J.D. Adams, S.C. Minne, and C.F. Quate, *Thin film characterization by atomic force microscopy at ultrasonic frequencies*, Appl. Phys. Lett, vol. 76, No. 14, Apr. 3, 2000, pp. 1950–1952.
- Jeffrey L. Hutter and John Bechhoefer, *Calibration of atomic-force microscope tips*, Rev. Sci. Instrum. 64(7), Jul. 1993, pp. 1868–1873.
- Mario B. Viani, et al., *Small cantilevers for force spectroscopy of single molecules*, Journal of Applied Physics, vol. 86, No. 4, Aug. 15, 1999, pp. 2258–2262.
- Constant A.J. Putman, et al., *Tapping mode atomic-force microscopy in liquid*, Appl. Phys. Lett. 64 (18), May 2, 1994, pp. 2454–2456.
- Marc P. Scherer, et al., *Experimental determination of the mechanical impedance of atomic force microscopy cantilevers in fluids up to 70 kHz*, Journal of Applied Physics, vol. 88, No. 5, Sep. 1, 2000, pp. 2912–2920.
- S.A. Elrod, et al., *Nozzleless droplet formation with focused acoustic beams*, J. Appl. Phys. 65(9), May 1, 1989, pp. 3441–3447.
- George T. Palocz, et al., *Rapid imaging of calcite crystal growth using atomic force microscopy with small cantilevers*, Appl. Phys. Lett, vol. 73, No. 12, Sep. 21, 1998, pp. 1658–1660.
- T. Sulchek et al., *High-speed atomic force microscopy in liquid*, Rev. Sci. Instrum., vol. 71, No. 5, May, 2000, pp. 2097–2099.
- T.J. Senden et al., *Experimental Determination of Spring Constants in Atomic Force Microscopy*, Langmuir, vol. 10, No. 4, (Feb. 10, 1994), pp. 1003–1004.
- W. Han et al., *A magnetically driven oscillating probe microscope for operation in liquids*, Appl. Phys. Lett. 69 (26), Dec. 23, 1996, pp. 4111–4113.
- B. Chu, et al., *Acoustic radiation pressure produced by a beam of sound*, J. Acoust. Soc. Am 72 (6), Dec. 1982), pp. 1673–1687.
- B. Hadimioglu et al., *High-Efficiency Fresnel Acoustic Lenses*, Ultrasonics Symposium, (1993), pp. 579–582.

* cited by examiner

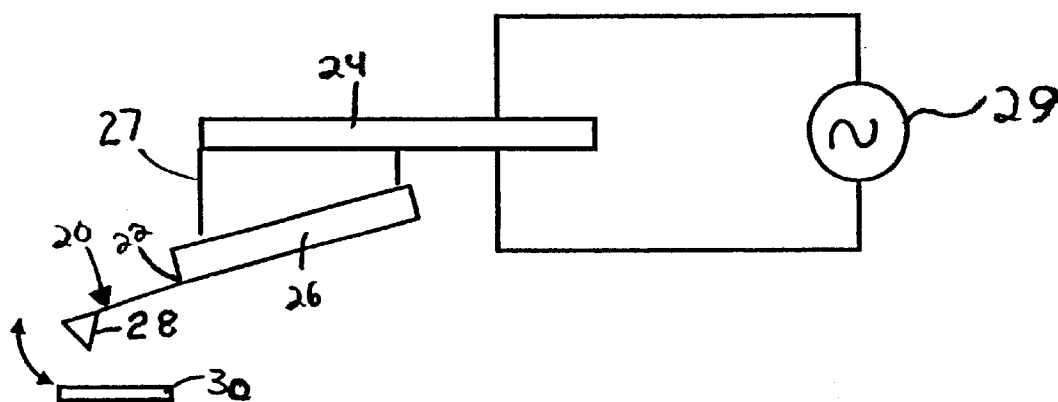


Figure 1A
(prior art)

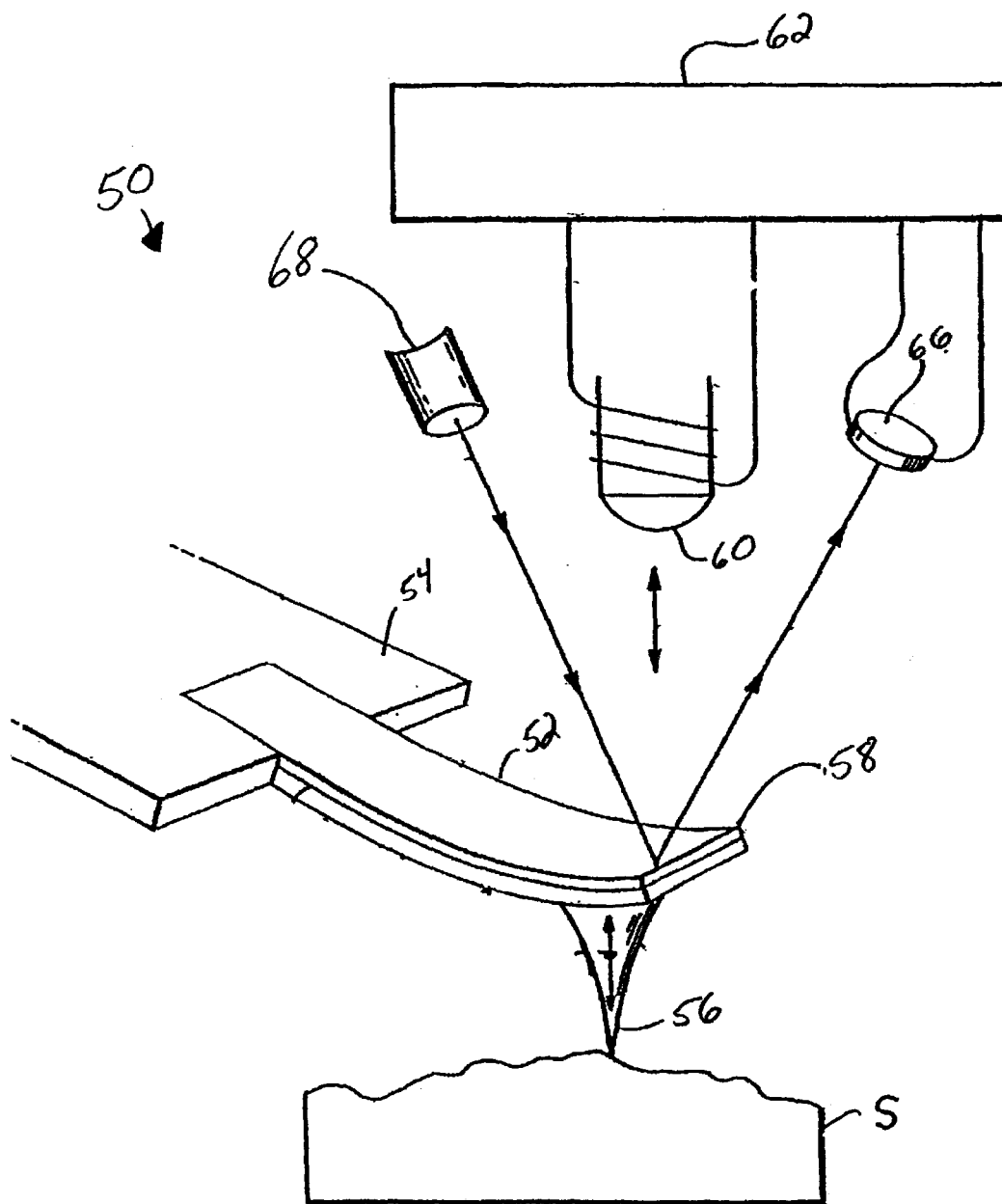


Figure 1B
(prior art)

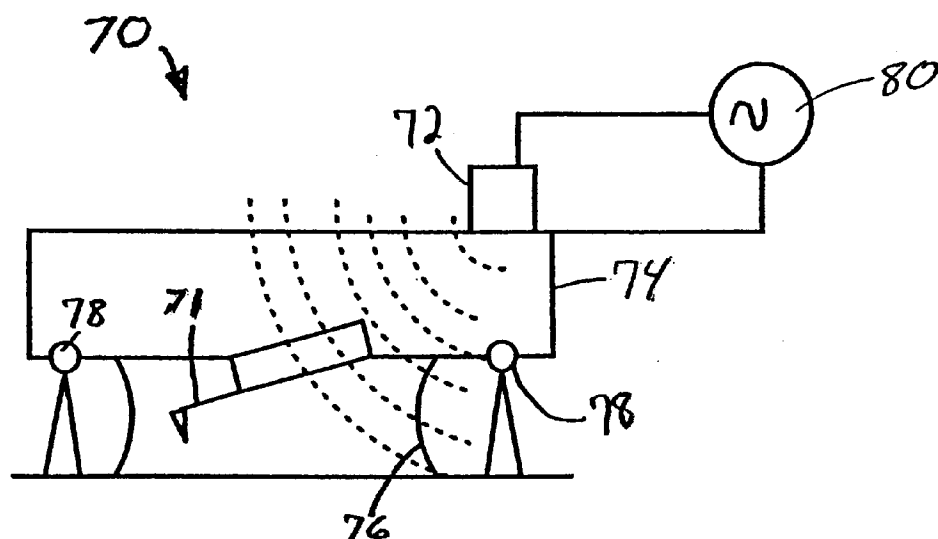


Figure 1C
(prior art)

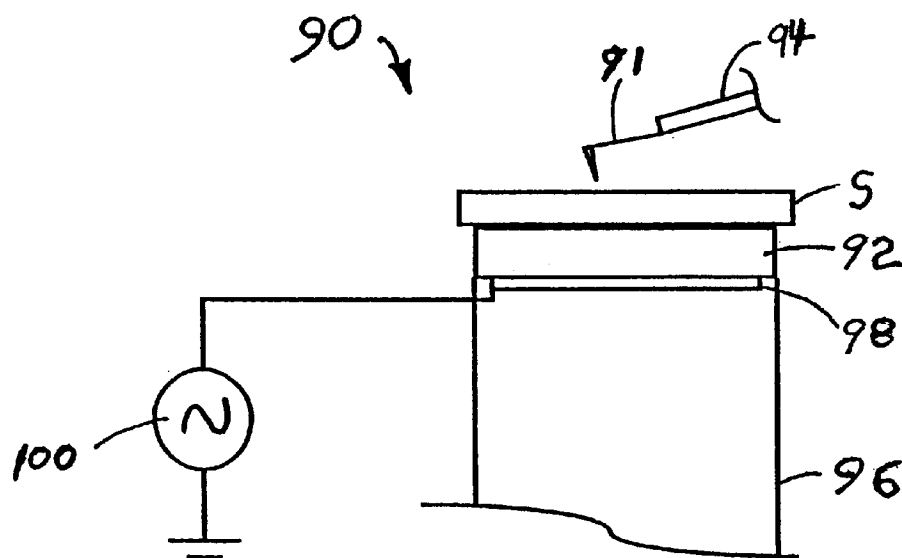


Figure 1D
(prior art)

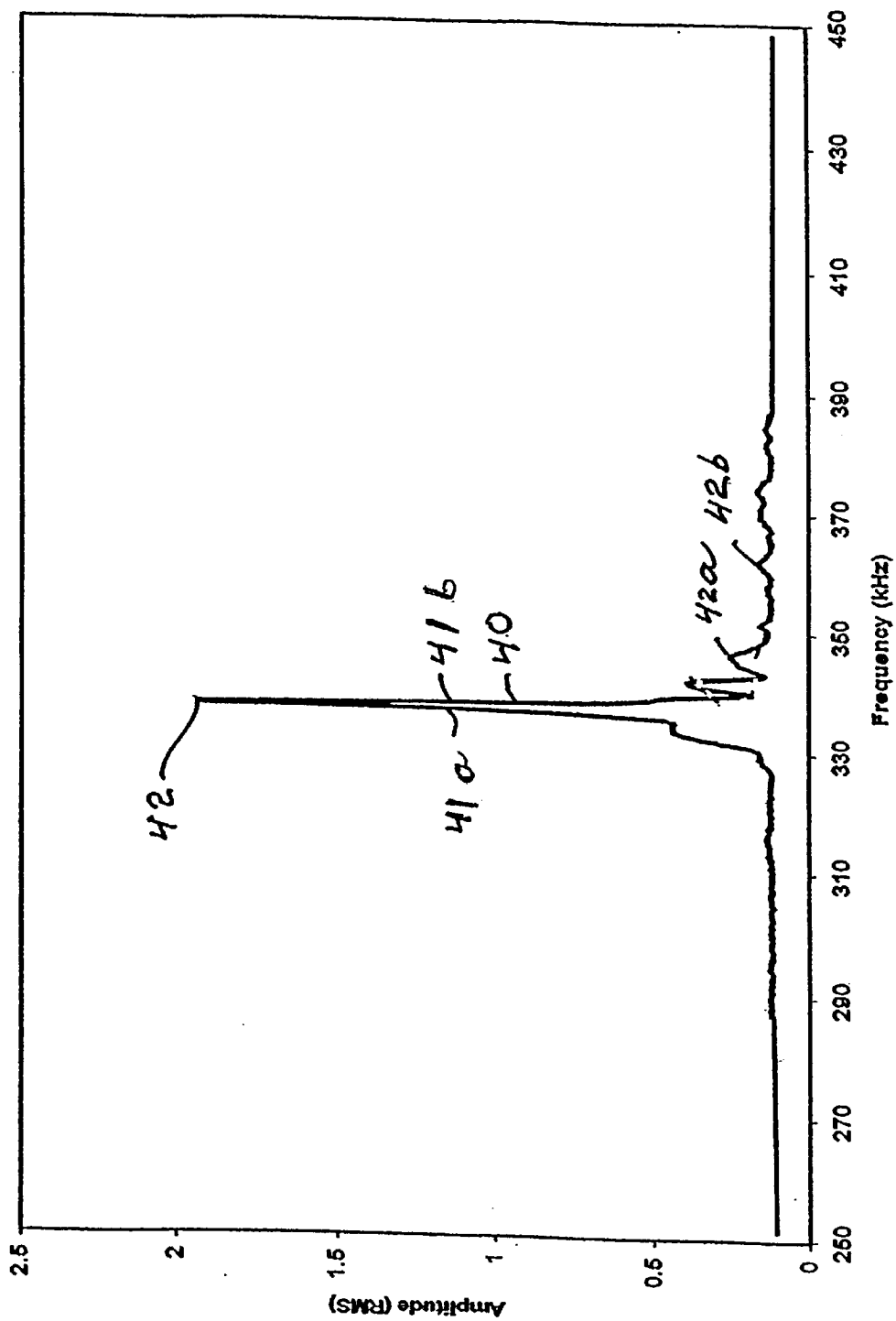


Figure 2A

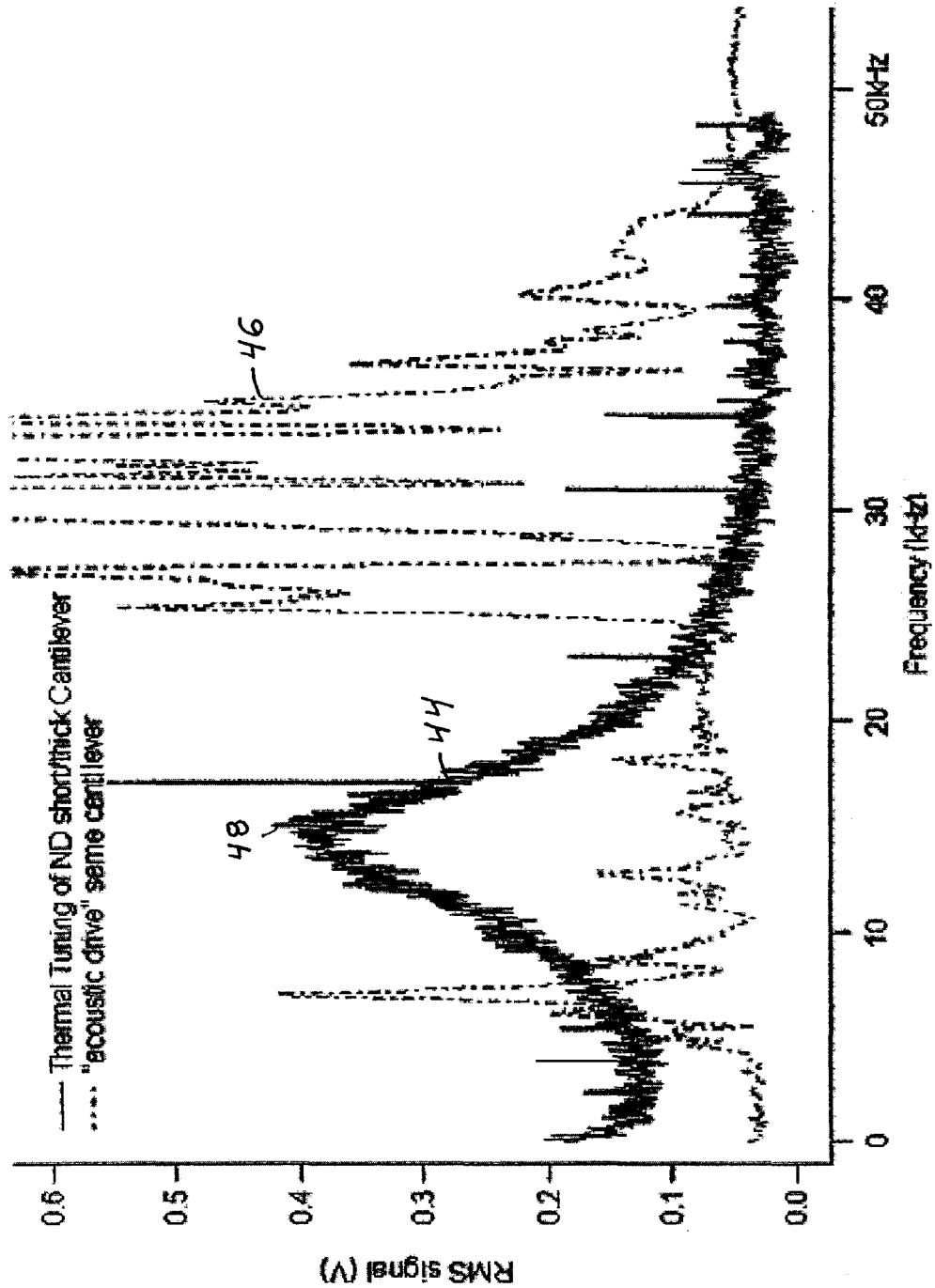


Figure 2B

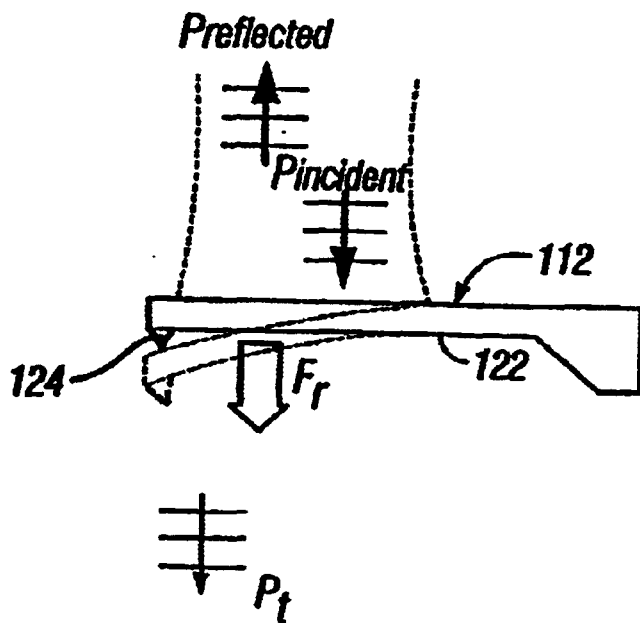


FIG. 3A

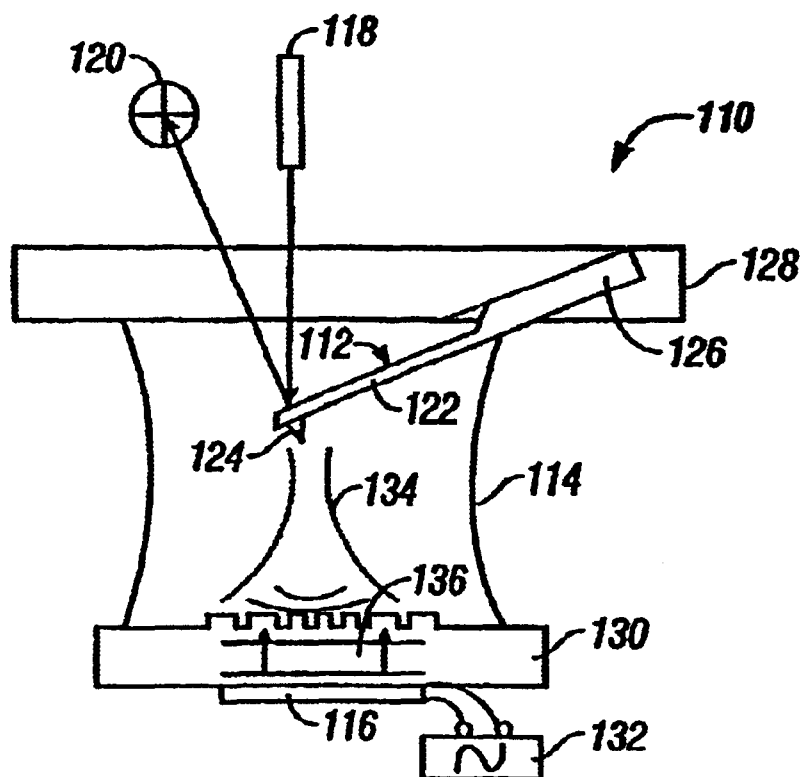


FIG. 3B

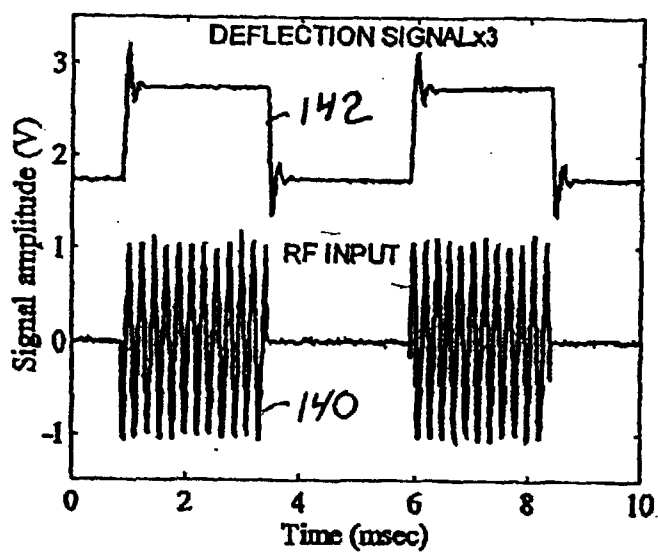


Figure 4

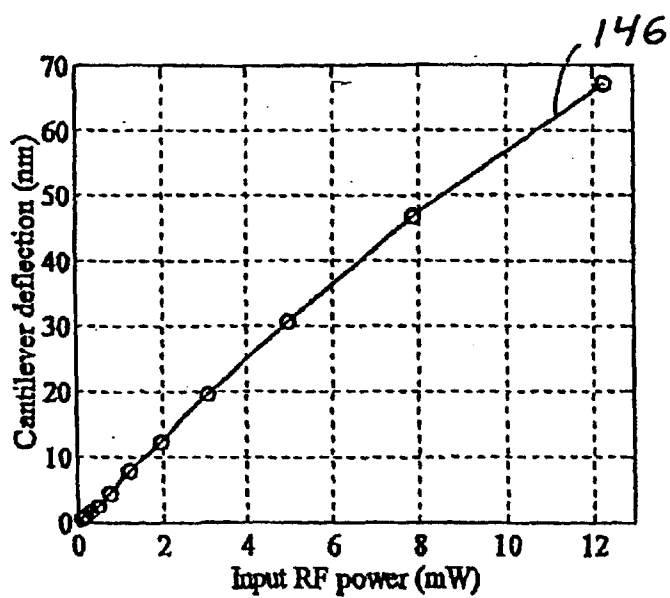


Figure 5

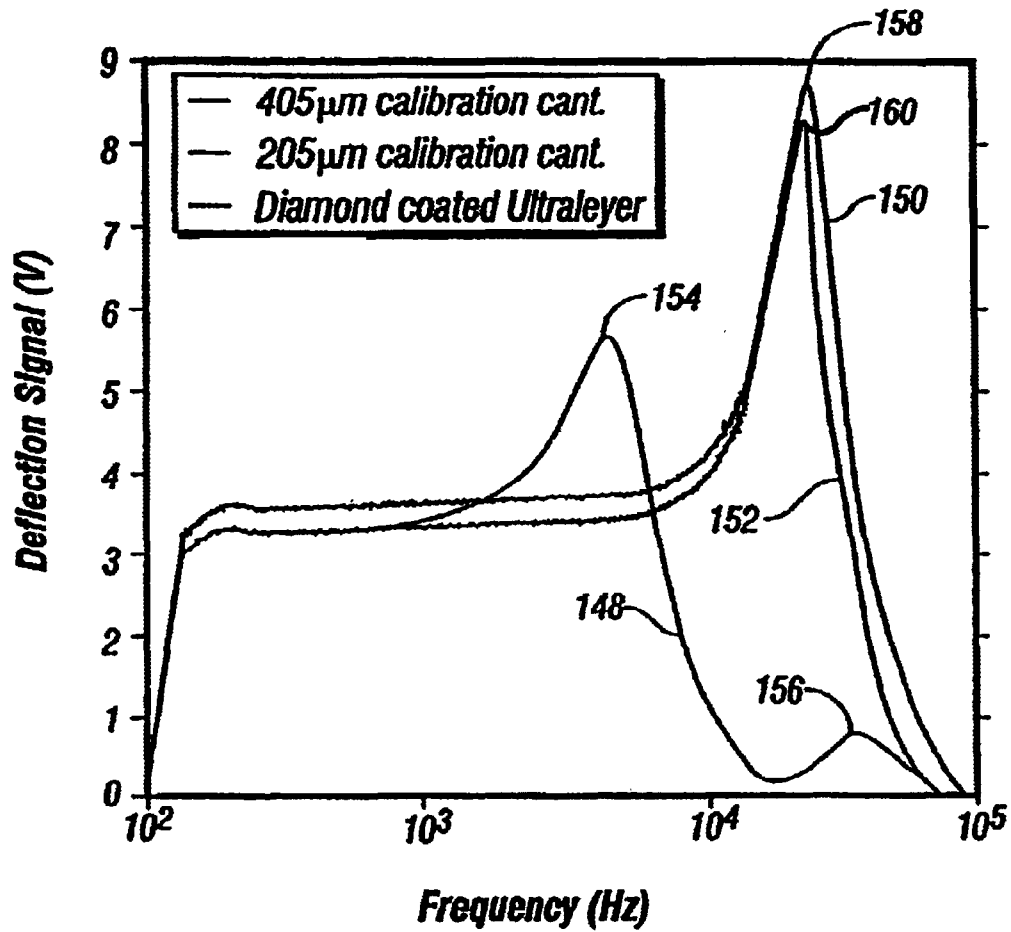


FIG. 6

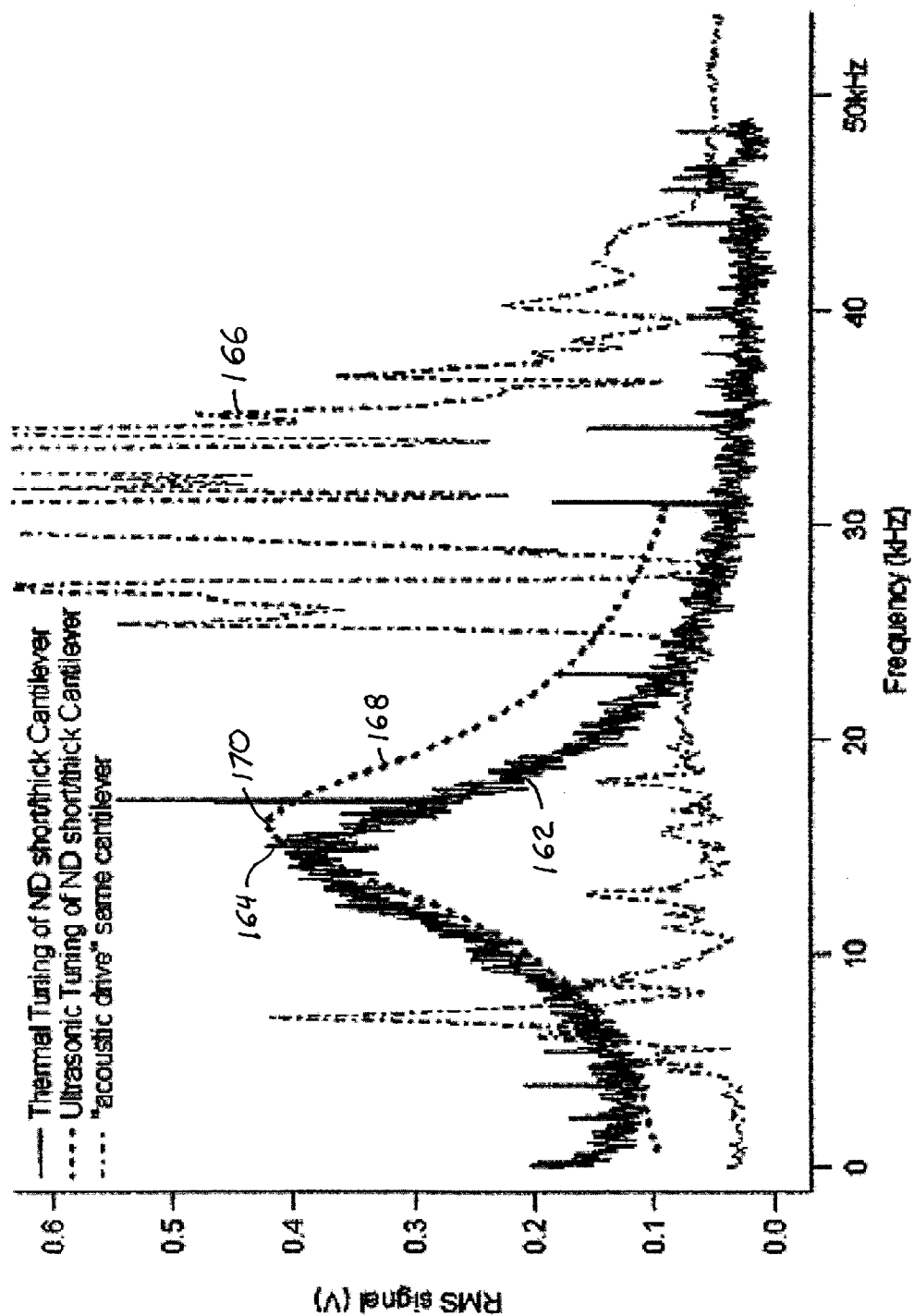


Figure 7A

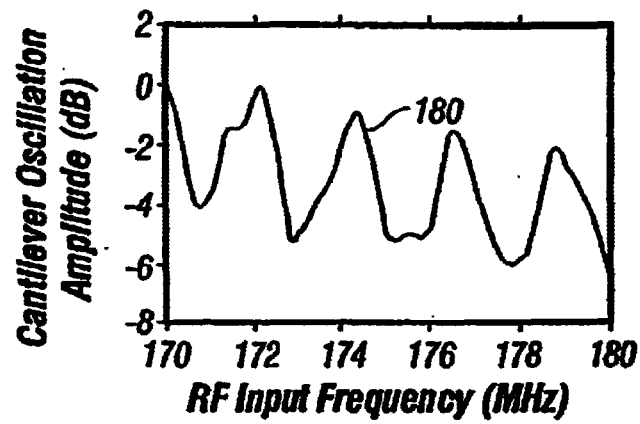


FIG. 7B

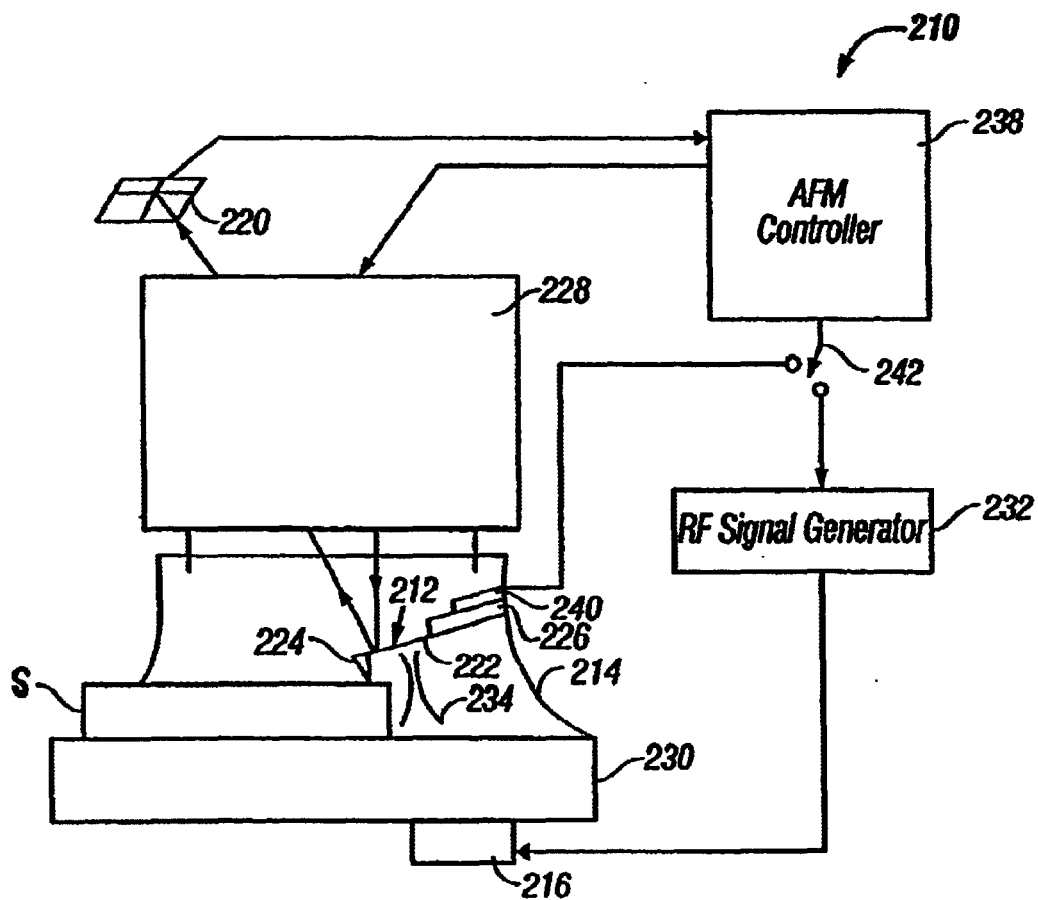


FIG. 8

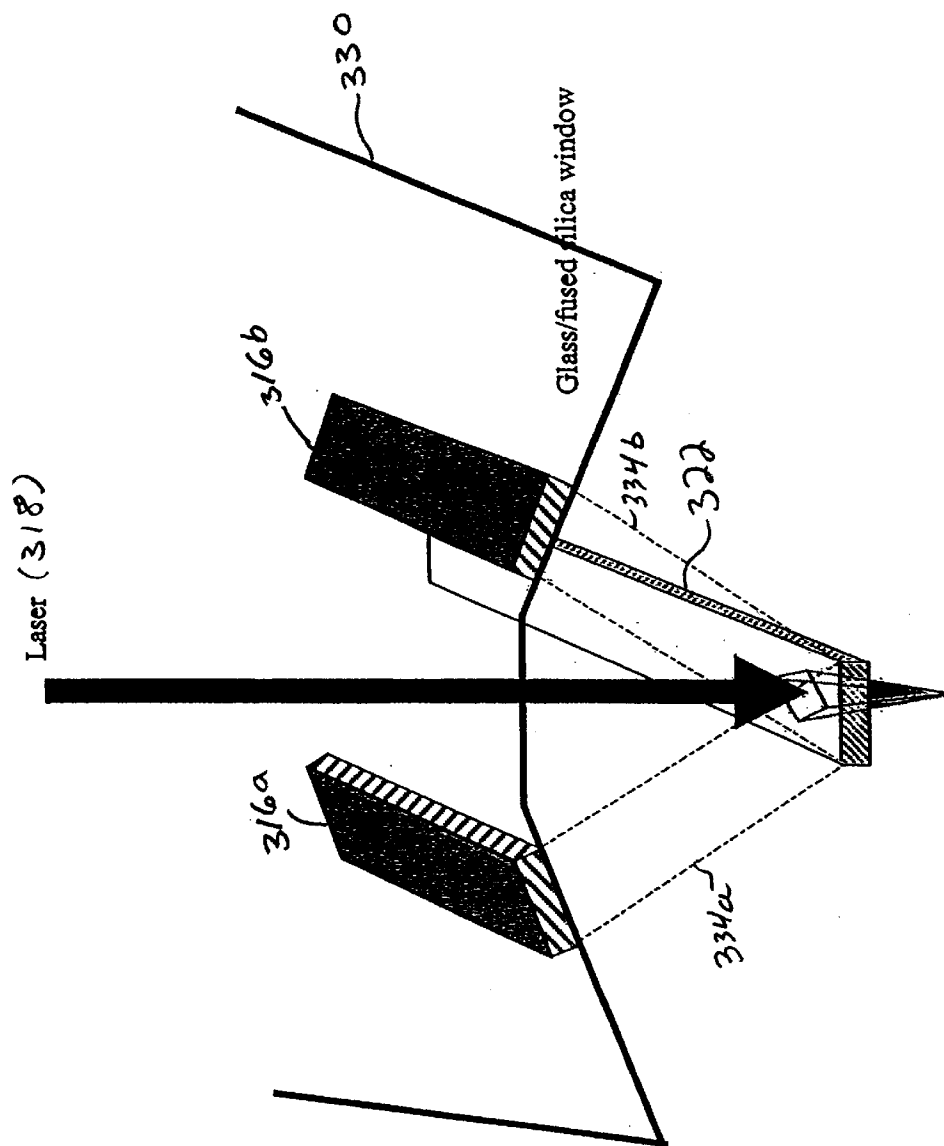
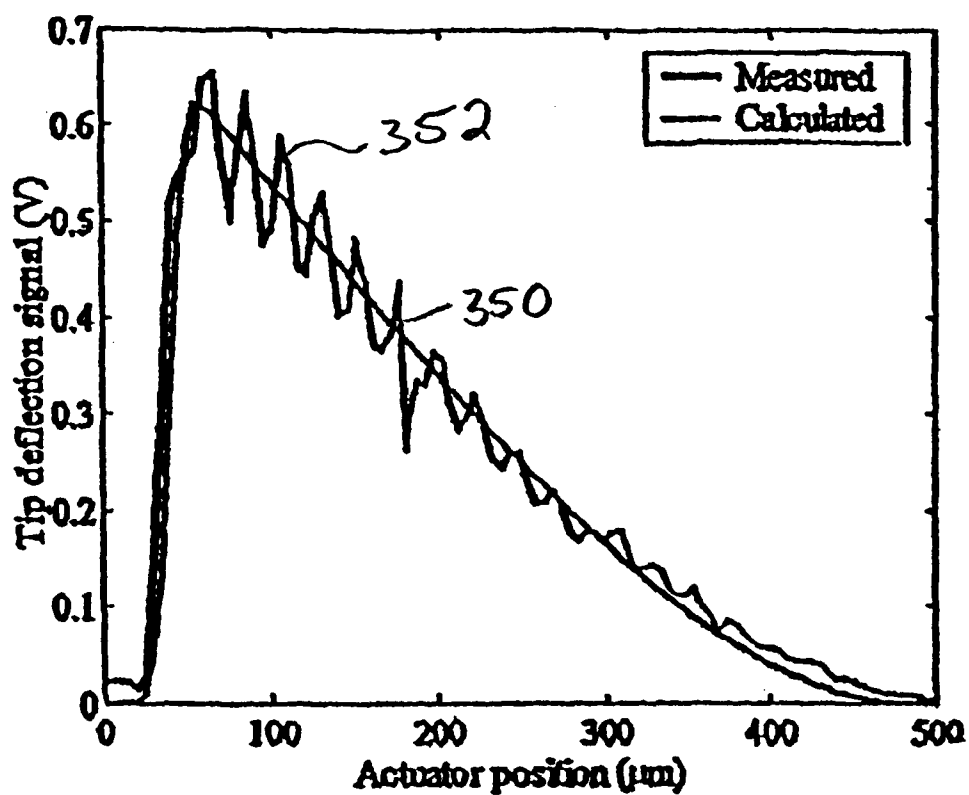
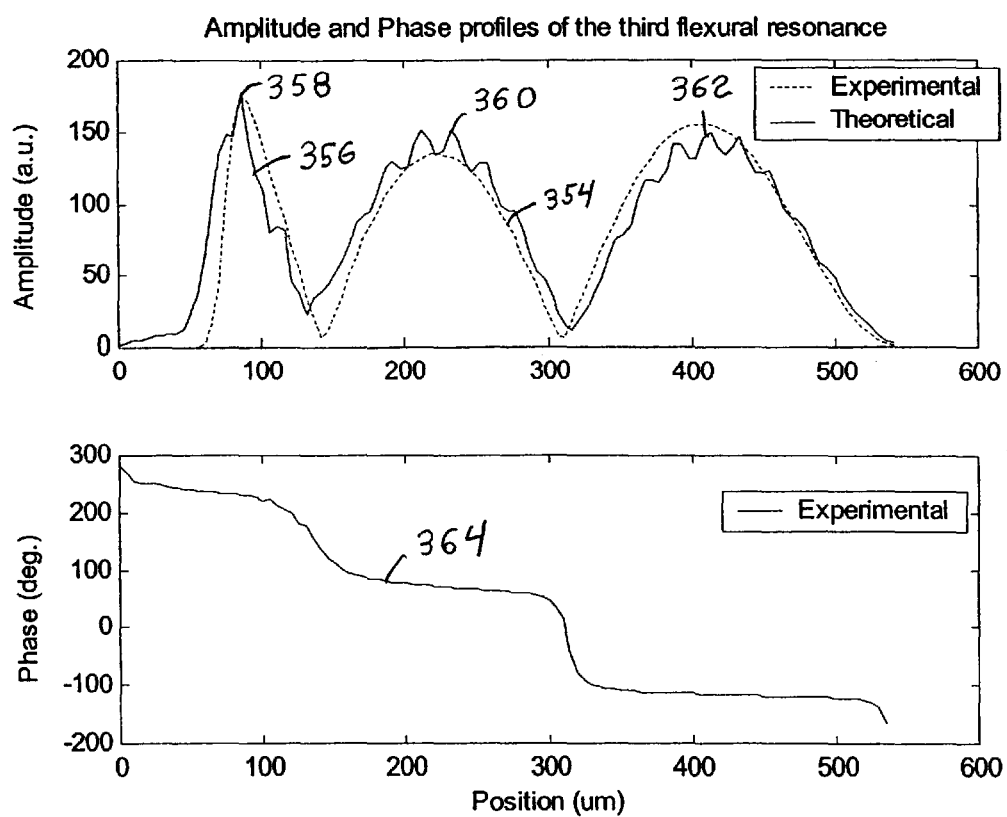


Figure 9

**Figure 10A**

**Figure 10B**

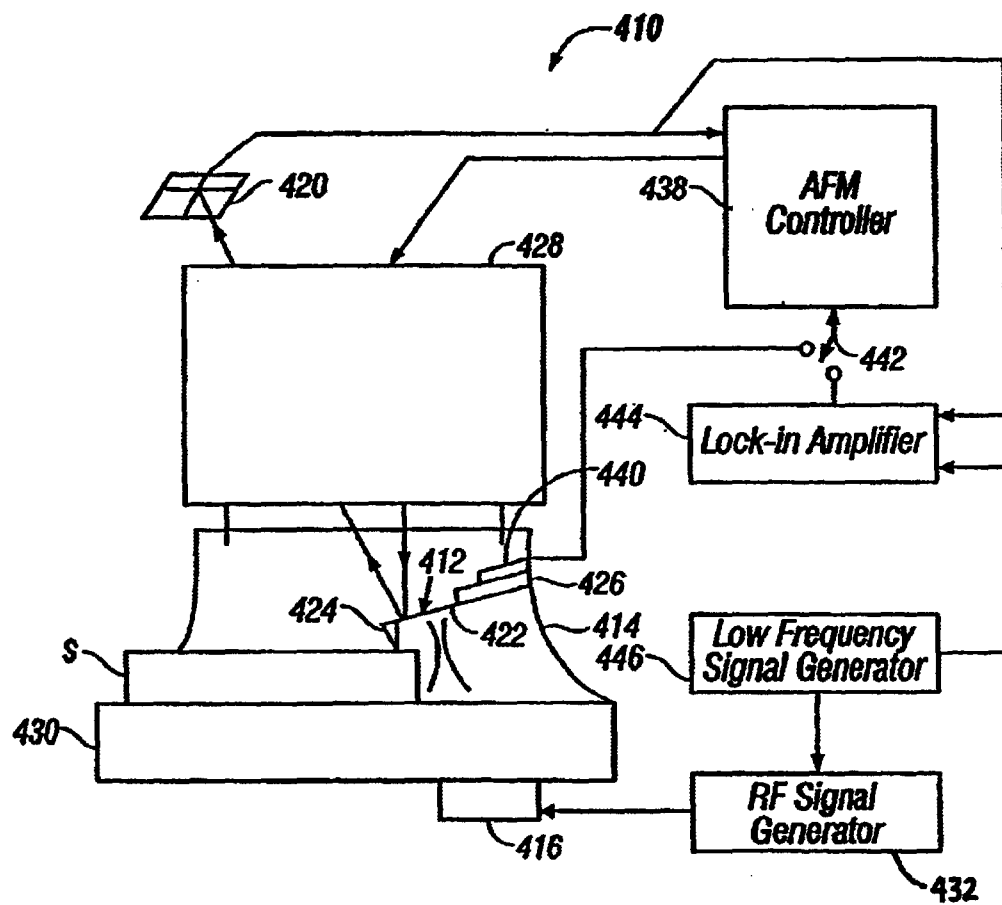
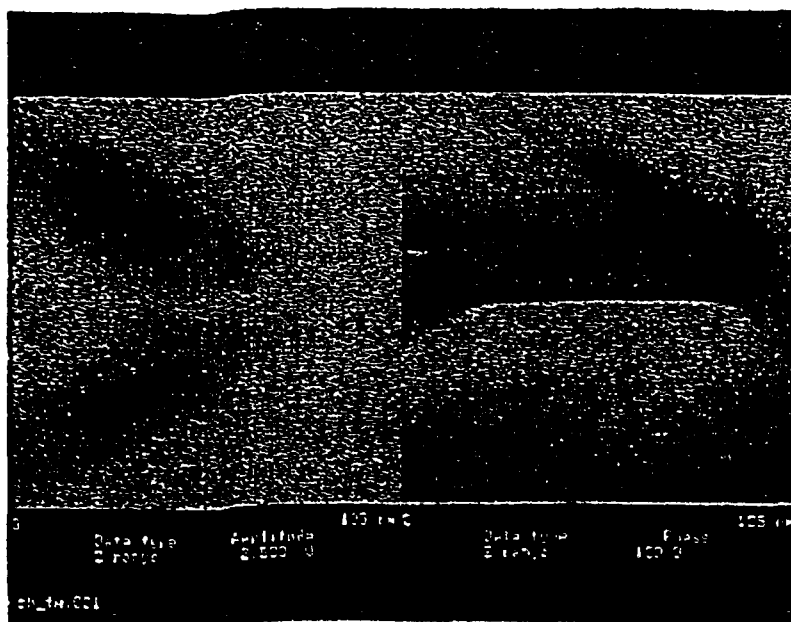
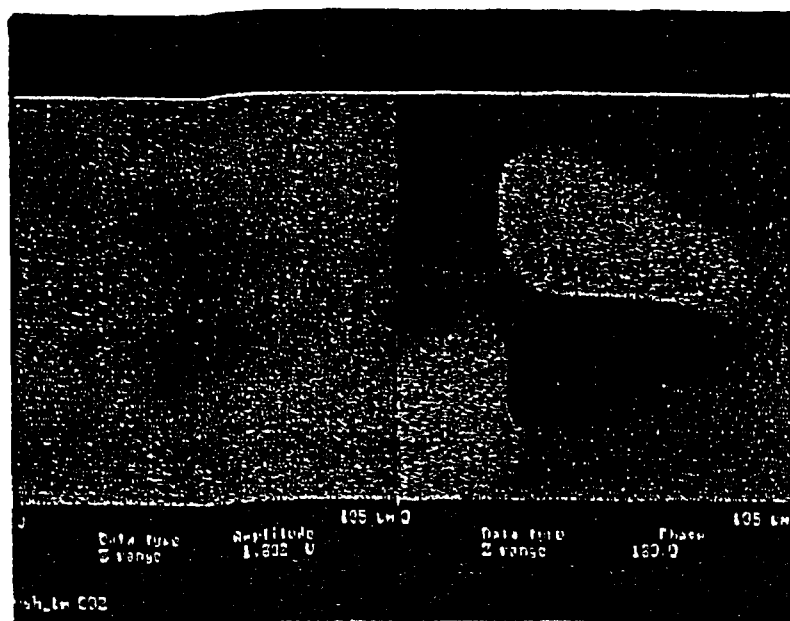
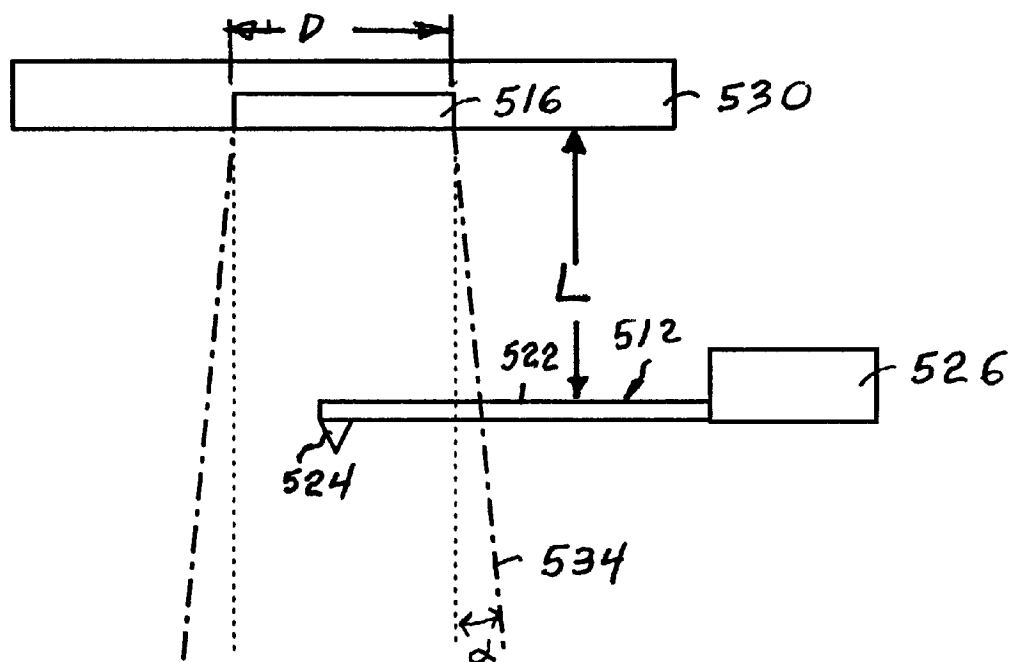
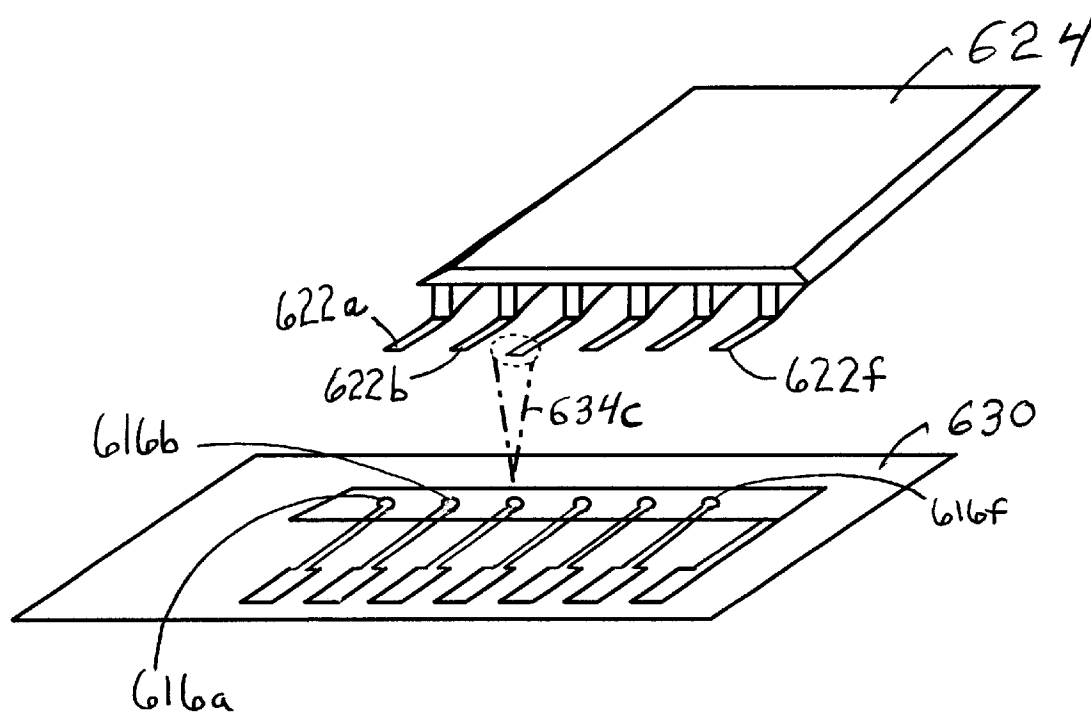
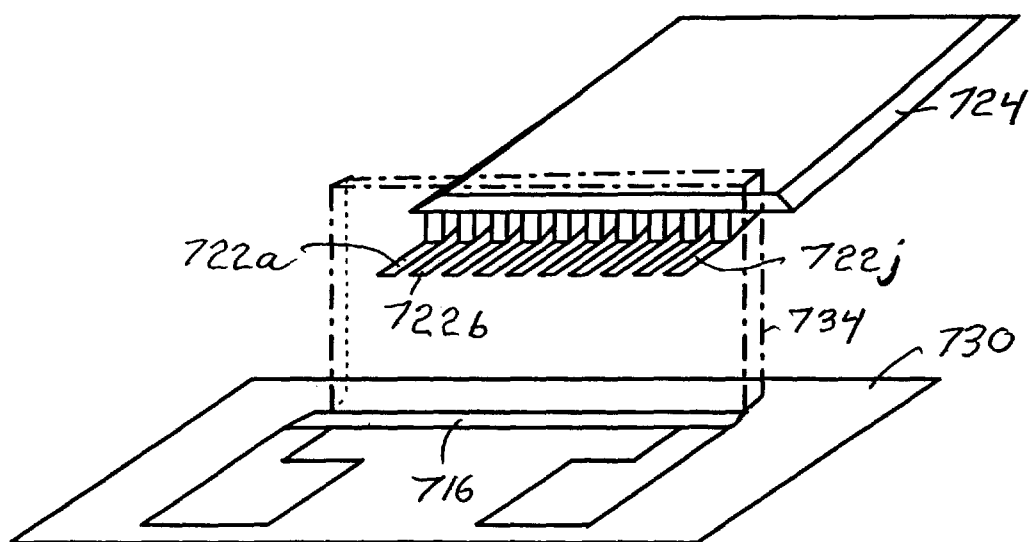


FIG. 11

**Figure 12A****Figure 12B**

**Figure 13**

**Figure 14A**

**Figure 14B**

METHOD AND APPARATUS FOR THE ULTRASONIC ACTUATION OF THE CANTILEVER OF A PROBE-BASED INSTRUMENT

CROSS-REFERENCE TO RELATED APPLICATION

Priority under 35 USC §1.119(e) is hereby claimed on prior U.S. Provisional Patent Application Serial No. 60/313, 911, filed Aug. 21, 2001, the subject matter of which is hereby incorporated by reference in its entirety.

BACKGROUND OF THE INVENTION

1. Field of the Invention

The present invention is directed to probe-based instruments and, more particularly, relates to a method and apparatus for driving a cantilever of such an instrument using acoustic radiation pressure generated by an ultrasonic actuator.

2. Description of Related Art

Several probe-based instruments monitor the interaction between a cantilever-based probe and a sample to obtain information concerning one or more characteristics of the sample. For example, scanning probe microscopes (SPMs) typically characterize the surface of a sample down to atomic dimensions by monitoring the interaction between the sample and a tip on the cantilever probe. By providing relative scanning movement between the tip and the sample, surface characteristic data can be acquired over a particular region of the sample, and a corresponding map of the sample can be generated.

The atomic force microscope (AFM) is a very popular type of SPM. The probe of the typical AFM includes a very small cantilever which is fixed to a support at its base and which has a sharp probe tip attached to the opposite, free end. The probe tip is brought very near to or into contact with a surface of a sample to be examined, and the deflection of the cantilever in response to the probe tip's interaction with the sample is measured with an extremely sensitive deflection detector, often an optical lever system such as described in Hansma et al. U.S. Pat. No. RE 34,489, or some other deflection detector such as strain gauges, capacitance sensors, etc. The probe is scanned over a surface using a high-resolution three axis scanner acting on the sample support and/or the probe. The instrument is thus capable of creating relative motion between the probe and the sample while measuring the topography or some other surface property of the sample as described, e.g., in Hansma et al. U.S. Pat. No. RE 34,489; Elings et al. U.S. Pat. No. 5,226,801; and Elings et al. U.S. Pat. No. 5,412,980.

AFMs may be designed to operate in a variety of modes, including contact mode and oscillating mode. In contact mode operation, the microscope typically scans the tip across the surface of the sample while keeping the force of the tip on the surface of the sample generally constant. This effect is accomplished by moving either the sample or the probe assembly vertically to the surface of the sample in response to sensed deflection of the cantilever as the probe is scanned horizontally across the surface. In this way, the data associated with this vertical motion can be stored and then used to construct an image of the sample surface corresponding to the sample characteristic being measured, e.g., surface topography. Alternatively, some AFMs can at least selectively operate in an oscillation mode of operation such as TappingMode™. In TappingMode™ the tip is oscil-

lated at or near a resonant frequency of the cantilever of the probe. The amplitude or phase of this oscillation is kept constant during scanning using feedback signals, which are generated in response to tip-sample interaction. As in contact mode, these feedback signals are then collected, stored, and used as data to characterize the sample.

Regardless of their mode of operation, AFMs can obtain resolution down to the atomic level on a wide variety of insulating or conductive surfaces in air, liquid or vacuum by using piezoelectric scanners, optical lever deflection detectors, and very small cantilevers fabricated using photolithographic techniques. Because of their resolution and versatility, AFMs are important measurement devices in many diverse fields ranging from semiconductor manufacturing to biological research.

One potentially problematic characteristic of AFMs and other probe-based instruments lies in the technique employed to provide an external force to deflect or oscillate the instrument's cantilever. In an AFM, the cantilever is typically oscillated using a piezoelectric drive, often known simply as a piezo drive. Referring to FIG. 1A by way of example in this type of system, the typical probe 20 includes a cantilever 22 that extends outwardly from a substrate 26 coupled to a piezoelectric drive 24 via a probe mount 27. Probe 20 also includes a tip 28 that is provided on the opposed, free end of the cantilever 20. The piezoelectric drive 24 can be selectively excited by a signal generator 29 to move the cantilever 22 up and down relative to a sample 30. When the instrument is configured for an oscillating mode of operation, the drive voltage is applied to the piezoelectric drive 24 to drive the cantilever 22 to oscillate at a frequency that is dependent upon the frequency of the drive voltage. This frequency is typically at or near the cantilever's resonant frequency, particularly when the instrument is operated in TappingMode™.

Such a traditional piezoelectric drive necessarily acts only on the base of the cantilever, not on the free end portion. It therefore must apply substantially greater forces to the cantilever to obtain a given deflection magnitude at the free end than it would if it were to apply forces directly to the free end or even to the body of the cantilever. This inefficiency limits the range of applications for this common type of piezo-electrically-driven probe.

For instance, the piezoelectric drive shown in FIG. 1A works well in air because the typical AFM cantilever can be excited to resonance in air relatively easily. This characteristic is quantified by the "quality factor" of a resonance of the cantilever. The quality factor, Q, denotes the sharpness of a cantilever's resonance curve as denoted by the ratio: $f_0/\Delta f$, where f_0 is the resonant frequency and Δf is the bandwidth between the half-power points of the curve as reflected by the half-peak amplitude points 41a and 41b on the curve 40 in FIG. 2A. The curve 40 demonstrates that the typical cantilever operating in air has a Q of 100–200 or even higher. The Q of a cantilever resonance is a measure of how much gain the resonance provides in an oscillating system. A resonance with a large Q can be excited to relatively large cantilever oscillation amplitudes with relatively small excitation forces. For operation in air or other gaseous environments, the cantilever typical piezoelectric drive usually has ample excitation force to drive the cantilever to produce a resonance peak 42 that is easily identified and distinguished from other, parasitic resonance peaks such as those of the mounts for the cantilever and the piezoelectric drive and the piezo drive itself (note the much smaller peaks 42a, 42b, etc. denoting these parasitic resonances).

Conversely, a cantilever operated in liquid such as water has a dramatically lower Q because the liquid damps the

oscillating cantilever. In fact, the typical cantilever operating in water has a Q of less than 30 and even less than 10. As a result, the typical piezoelectric drive does not have enough gain to excite the cantilever sufficiently to produce a resonance peak that is easily located and differentiated from parasitic resonances. This effect is discussed below in conjunction with FIG. 2B.

Specialized cantilever drives are available that act along the length of the cantilever rather than only on the base. One such drive is the so-called magnetic drive. Referring to FIG. 1B, the typical magnetic drive system **50** has a magnetic cantilever **52** that is driven by an electromagnetic drive. The cantilever **52** has a fixed base rigidly attached to a support **54** and bears a tip **56** on its free end that interacts with sample S. The cantilever is also rendered magnetic by coating one or more of its surfaces with a magnetic layer **58**. The electromagnetic drive comprises at least one electromagnet coil **60** spaced from the cantilever **52**. The coil(s) can be energized by a controller **62** including a signal generator to impose a variable magnetic field on the magnetic layer **58**. The magnetic field produces a torque on the cantilever **52** of a magnitude that increases with the amplitude of the applied magnetic field acting on the layer **58**. By controlling the amplitude of the applied magnetic field, the cantilever **52** can be deflected as desired while the tip **56** interacts with the sample S. This deflection is monitored by a photodetector **66** receiving reflected light transmitted by a laser **68**. In the usual case in which the magnetic drive **60** is controlled to maintain a specified characteristic of cantilever deflection constant during scanning, an output signal related to the amplitude of the signal provides an indication of a surface force applied to the probe. A magnetic drive system having these characteristics is described in greater detail in U.S. Pat. No. 5,670,712 to Cleveland, the subject matter of which is hereby incorporated by reference by way of background.

A magnetic drive system has inherent limitations that considerably restrict its range of applications. For instance, it requires a special magnetically coated cantilever and, accordingly, cannot be used in applications in which the cantilever is not capable of being coated with a magnetic material. It also is not usable in applications in which magnetic properties of the sample and/or the environment cause unwanted deflection of the cantilever and produce errors into the measurements. The practical operating ranges of the magnetic drive system are also limited. A typical magnetic drive coil may operate with a current exceeding an amp and result in a cantilever deflection on the order of 1–100 nm at the cantilever resonance frequency. Even at this coil current, the heat load generated can cause thermal drift errors in the measurement of the AFM. The frequency range of the magnetic drive system is also limited by the inductance of the drive coil. Higher actuation forces can be achieved by using more loops in the drive coil, but this also increases the inductance and limits the maximum operating frequency. With the limits of inductance and maximum heat load, the typical magnetic drive operates with less than 50 kHz and with oscillation amplitudes of less than 30 nm. For example, the MAC-Mode™ magnetic drive system, sold by Molecular Imaging, advertises an operating range of 5–30 kHz and a maximum amplitude of 30 nm.

Another instrument having a cantilever driven remotely from its base utilizes the so-called acoustic drive. Referring to FIG. 1C, in an instrument **70** of this type, a cantilever **71** and a piezoelectric drive **72** are mounted on a common head **74** in a spaced-apart relationship. The head **74** is mounted above a fluid cell **76** by mounting balls **78** or other supports so that the cantilever **71** extends into the fluid cell **76** so as

to interact with a sample (not shown) in the cell. The piezoelectric drive **72** can be excited by a signal generator **80** to generate acoustic waves that propagate through the glass walls of the fluid cell **76**, through the fluid in the cell **76**, and onto the cantilever **71**, causing the cantilever **71** to oscillate. An acoustic drive having these characteristics is disclosed, for example, in Putman et al in “Tapping Mode Atomic Force Microscopy in Liquids” Applied Physics Letters 64: 2454–2456.

Acoustic drive has distinct disadvantages that limit its effectiveness. For instance, the acoustic energy also impinges on many other components of the system, such as mounts for the cantilever and the piezoelectric drive, the fluid cell, and even the fluid exciting, resonances in those components. These resonances can be difficult to distinguish from the cantilever resonance. The acoustic drive also has sufficient actuation force at a limited selection of operation frequencies and it can be a challenge to match the cantilever resonance with the operation frequency of the acoustic actuator. If a user selects a resonance that does not overlap with the cantilever resonance, the measurements may be unstable.

An ultrasonic force microscope (UFM) is a scanning probe microscope that uses high frequency acoustic waves to image the mechanical properties of a sample, often showing sub-surface contrast. Specifically, referring to FIG. 1D, a UFM **90** includes a cantilever **91** having a base fixed to a stationary support **94**, a sample support **92** located beneath the cantilever **91**, and on an XYZ scanner **96** that supports the sample support. An ultrasonic actuator **98** such as commercial ultrasound transducer mounted on the bottom of the sample support **92** and is excited by an RF voltage from an RF signal generator **100**. The ultrasonic actuator **98** is relatively large (typically a centimeter or more in diameter) with a resonant frequency often in the low-MHz range. When it is excited by the RF signal generator **100**, it generates ultrasonic waves that impinge over a broad area of the sample S. Some of the incident ultrasonic energy is reflected or absorbed, and some penetrates the sample S and then impinges on the cantilever **91**, causing the cantilever **91** to deflect away from the sample surface. The magnitude of the cantilever deflection is related to the percentage of the energy that penetrates the sample S and, accordingly, the, reflects variation in sample properties such as density. Accordingly, as the sample S is scanned relative to the probe using the scanner **96**, variations in cantilever deflection can be detected to provide information concerning the sample. In addition, while UFM's have been in use for almost a decade, no one has adapted an ultrasonic device as a general purpose cantilever actuator capable of deflecting the cantilever at a wide range of frequencies.

Turning to FIG. 2B mentioned above, the plots demonstrate the frequency response of a typical AFM cantilever to excitation. The curve **44** plots the actual or true response of a relatively short and thick cantilever in water as determined by a known process called a “thermal tune.” A thermal tune measures the natural intrinsic motion of the cantilever in response to the temperature of its surroundings. Basically, the “heat bath” that surrounds the cantilever provides the energy to naturally oscillate at a very small amplitude, usually sub-nm. Since the cantilever oscillation amplitude due to the thermal energy is so small, thermal tunes cannot be used for image data acquisition, but they do provide a very clean representation of the true oscillatory response of the cantilever. The curve **46** plots the detected response of the same cantilever as it is driven acoustically by a piezoelectric drive (FIG. 1C). The true response as denoted in

curve 44 has a sharp peak 48 at the fundamental resonance of about 15 kHz. However, when the cantilever is driven acoustically by a piezoelectric drive, fluid damping and other effects reduce that response to the point that the cantilever resonance peak cannot be differentiated from parasitic resonance peaks.

Hence, the need has arisen to provide a probe-based instrument that has an actuator that drives the cantilever 50 as to produce a "clean" frequency response, preferably by driving the cantilever body rather than the base, but that is versatile in bandwidth/or types of measurements.

The need has also arisen to provide an improved method of driving a cantilever of a probe-based instrument.

SUMMARY OF THE INVENTION

In accordance with a first aspect of the invention, one or more the above-identified needs is met by providing a probe-based instrument having a cantilever that is deflected by directing acoustic waves onto the body of the cantilever rather than by moving the base of the cantilever. The cantilever is deflected by a second order force, also known as an acoustic radiation force, generated by beams of ultrasonic energy generated by an ultrasonic actuator such as a zinc oxide transducer. The ultrasonic actuator is supplied with an oscillating RF voltage that may be continuous or varied in a quasistatic manner to apply a constant or changing force to the cantilever. The RF voltage may also be modulated at any frequency from DC to many MHz, thus providing an ideal drive force for oscillating the cantilever over an extremely wide range of frequencies. Driving the body of the cantilever with an ultrasonic actuator produces a much higher localized force than can be achieved through the use of a traditional piezoelectric actuator and, accordingly, permits a "clean" frequency response where the resonance peak is easily identified and differentiated from parasitic resonance peaks. This, in turn, dramatically improves the accuracy, precision, and stability of the measurement, and increases the system's bandwidth, particularly when the cantilever operates in a liquid. The method implemented by the invention can be used to actuate cantilevers with arbitrary shapes and materials, eliminating the requirement for magnetic or piezoelectric coatings on the cantilever. The method and system of the preferred embodiments also are useful in imaging in liquids and quantitative measurements of surfaces and molecular-scale samples in liquids. The method and system of the preferred embodiments also are useful in AFM measurements in other fluids including air.

The improved frequency response of the ultrasonic actuator of the preferred embodiment also yields a dramatically higher bandwidth than traditional piezoelectric actuators, rendering them useful in a variety of applications and with a variety of cantilevers beyond those available with conventional piezoelectric actuators.

The beam is preferably "shaped", i.e., manipulated to limit unwanted propagation in directions other than toward the cantilever, so that ultrasonic energy impinges at least primarily on the cantilever. Two suitable techniques for shaping the beam are focusing and collimation. Ultrasonic beams can be focused on the cantilever using a Fresnel lens or another focusing device located between the ultrasonic actuator and the cantilever. Collimation requires only that the ultrasonic actuator be suitably sized, positioned, and driven to reduce beam divergence sufficiently to achieve the desired effect.

Cantilever deflection may be measured by a conventional photodetector, in which case the photodetector, a laser, and

the ultrasonic actuator are all preferably positioned on a common side of the cantilever opposite the sample support. Cantilever deflection may also be detected using another device such as a simple interferometer located over the cantilever body.

These and other features and advantages of the invention will become apparent to those skilled in the art from the following detailed description and the accompanying drawings. It should be understood, however, that the detailed description and specific examples, while indicating preferred embodiments of the present invention, are given by way of illustration and not of limitation. Many changes and modifications may be made within the scope of the present invention without departing from the spirit thereof, and the invention includes all such modifications

BRIEF DESCRIPTION OF THE DRAWINGS

A preferred exemplary embodiment of the invention is illustrated in the accompanying drawings in which like reference numerals represent like parts throughout, and in which:

FIG. 1A is a schematic view of a conventional AFM having a piezoelectric drive, appropriately labeled "prior art";

FIG. 1B is a schematic view of a conventional magnetic drive system, appropriately labeled "prior art";

FIG. 1C is a schematic view of a conventional acoustically driven microscope, appropriately labeled "prior art";

FIG. 1D is a schematic view of a conventional ultrasonic force microscope (UFM), appropriately labeled "prior art";

FIG. 2A is a graph illustrating the response of a typical AFM to an oscillating excitation when the cantilever is operating in air;

FIG. 2B is a response of a typical AFM cantilever to oscillating excitation when the cantilever is operating in a liquid;

FIG. 3A schematically illustrates the deflection of a cantilever of an AFM using ARP;

FIG. 3B schematically illustrates an ultrasonically actuated AFM constructed in accordance with a first embodiment of the present invention and incorporating the ARP deflected cantilever of FIG. 3A and an associated ultrasonic actuator;

FIG. 4 is a graph illustrating signal amplitude versus time of both drive input and measured cantilever deflection of the instrument of FIG. 3B;

FIG. 5 is a graph of cantilever deflection versus input power for the instrument of FIG. 3B;

FIG. 6 is a graph of deflection versus frequency of various cantilevers usable in the system of FIG. 3B;

FIG. 7A is a family of curves illustrating the response of a cantilever of the system of FIG. 3B to an oscillating input drive signal in liquid using both a prior art actuator and the actuator of FIG. 3B;

FIG. 7B is a graph illustrating the response of the cantilever of the system of FIG. 3B to changes in RF input frequency;

FIG. 8 schematically illustrates an ultrasonically actuated AFM constructed in accordance with a second embodiment of the invention in which the AFM is configured for TappingMode operation;

FIG. 9 is a schematic top plan view of the ultrasonic drive for an AFM constructed in accordance with a third embodiment of the invention in which the ultrasonic actuator and

detector are positioned on a common side of the cantilever opposed to the sample supports;

FIG. 10A is a graph of probe tip deflection versus beam impingement position for low frequency excitation of the cantilever of the instrument of FIG. 8;

FIG. 10B is a graph of probe tip deflection versus beam impingement position for high frequency excitation of the cantilever of the instrument of FIG. 8;

FIG. 11 schematically illustrates an ultrasonically actuated AFM constructed in accordance with a fourth embodiment of the invention in which the AFM is configured to take elasticity measurements;

FIG. 12A is a comparative array of amplitude and phase images obtained by operating the AFM of FIG. 11 in a first torsional mode;

FIG. 12B is a comparative array of amplitude and phase images obtained by operating the AFM of FIG. 11 in a second torsional mode;

FIG. 13 schematically illustrates an ultrasonically actuated AFM constructed in accordance with another embodiment of the invention in which the ultrasonic actuator and related components are configured to transmit a collimated or minimally divergent beam onto the cantilever as opposed to a focused beam;

FIG. 14A schematically illustrates a portion of an ultrasonically actuated cantilever array showing an array of ultrasonic actuators matched to an array of cantilevers; and

FIG. 14B schematically illustrates a portion of an ultrasonically actuated cantilever array with a single rectangular ultrasonic actuator arranged to excite the array of cantilevers.

DETAILED DESCRIPTION OF THE PREFERRED EMBODIMENT

As discussed briefly in the Summary section above, the invention lies in the use of ultrasonically generated acoustic radiation pressure to deflect a cantilever of a probe-based instrument. It is particularly well suited, but in no way limited, for instruments configured to take measurements in a liquid such as water or other aqueous solutions. A variety of SPMs and other instruments may benefit from this technique. Hence, while several different AFMs incorporating the invention will now be described by way of example, it must be emphasized that the invention is not limited to the described embodiments or even to AFMs in general. To the contrary, it is applicable to virtually any probe-based instrument in which a cantilever is deflected by directing a beam of ultrasonic energy at the cantilever to apply ultrasonically generated acoustic radiation pressure to the cantilever. The beam preferably is shaped, either by focusing substantially onto a surface of the cantilever, or by generating a sufficiently collimated or minimally divergent beam to permit to at least a portion of the beam to strike the cantilever. A variety of different ultrasonic actuators and associated drives may be employed to achieve these effects, and cantilever deflection may be measured using a variety of techniques.

Theory of Operation

A plane target placed in the path of an acoustic wave beam in an unconfined medium experiences a time averaged force per unit area. This pressure is known as the "Langevin acoustic radiation pressure" (ARP). The invention lies in the use of the forces imposed by the application of this pressure to deflect the cantilever of a probe-based instrument. One such probe 112 is illustrated schematically in FIG. 3A. The

probe 112 includes a cantilever 122 and a tip 124 mounted or otherwise provided on a free end of the cantilever 122. The ARP is related to the average energy density, U , at the target surface and it may if desired be localized at a specific location on the cantilever 122 by placing that location of the cantilever 122 at the focal plane of an acoustic lens. As a simple model based on plane waves, it can be assumed that, at the focal plane, a time harmonic acoustic pressure wave of amplitude, P_i , is normally incident on a cantilever immersed in a liquid and that the wave is reflected with a complex pressure reflection coefficient, F , at an angular frequency, $\omega=2\pi f$. This reflection coefficient can be considered as a weighted average over the incident spectrum of plane waves that would be included in a focused beam. In this case, the time-averaged energy density at the cantilever surface will be given by

$$U = \frac{P_i^2}{2\rho c^2} (2|\Gamma|^2) \quad \text{Equation (1)}$$

where ρ is the bulk density of fluid, c is the speed of sound in the liquid and $|\Gamma|$ denotes the absolute value of the reflection coefficient. Using the relation that the average intensity of the incident beam is given by $I_i=P_i^2/(2\rho c)$, the Langevin radiation pressure on the cantilever 122, Ω , can be expressed in terms of the intensity as

$$\Omega = \frac{I_i}{c} (2|\Gamma|^2) \quad \text{Equation (2)}$$

The total force applied to the cantilever in the direction of propagation of the incident wave can be found by integrating the radiation pressure. Accordingly, the total applied force is proportional to the average power incident on the cantilever 122. Note that the discussion above neglects the absorption of the ultrasonic energy in the beam and in the fluid medium. In case of absorption in the fluid medium, acoustic streaming can be induced. The fluid flow induced by this mechanism can generate additional forces on the cantilever. The losses in the cantilever are generally very small and hence can be neglected. Also note that Equations 1 and 2 apply to cantilever actuation applications in air. Since the velocity of sound in air is approximately 330 m/s and $|\Gamma|\sim 1$, the same amount of force can be applied to the cantilever with $1/5^{th}$ of the acoustic power. The high attenuation of ultrasonic waves in air may limit the frequency of operation.

The localization of the radiation force can be estimated using the relations for diffraction limited focused acoustic beams. For an acoustic lens with an F-number equal to one, the 3-dB diameter, d_s , of the diffraction limited beam at the focal plane is given by the relation $d_s \approx 1.02\lambda$, where λ is the wavelength of the time harmonic acoustic wave in the fluid. For example, in water ($c \approx 1.5 \times 10^3$ m/s), the diameter of the beam at the focal plane will be between $5\text{ }\mu\text{m}$ to $10\text{ }\mu\text{m}$ when the RF drive frequency, f , is in the 150–300 MHz range. According to Equation 2, for 150 μW incident average acoustic power at the focal plane of the lens, $P_i = \pi d_s^2 I_i / 4$, the force applied to the AFM cantilever will be 145 nN, assuming perfect reflection at the water/cantilever interface ($|\Gamma|=1$). Both these frequency and power levels are typically used for acoustic microscopy and acoustic ink printing applications.

It has been discovered that two separate high frequency beams with slightly different frequencies can be used to generate radiation pressure at the difference frequency to generate elasticity images. The acoustic radiation pressure

field can be localized by intersecting the high frequency beams at the desired location. This results in an effective amplitude modulation. It has also been discovered that modulated sonic beams can be used to generate acoustic forces to image mechanical properties of a variety of objects with high spatial resolution. These uses are discussed, for example, in: M. Fatemi and J. F. Greenleaf, "Ultrasound-stimulated vibro-acoustic spectrography," *Science*, 280, pp. 82-5, 1998; U.S. Pat. No. 5,991,239, Confocal Acoustic Force Generator; U.S. Pat. No. 5,921,928, Acoustic Force Generation by Amplitude Modulating a Sonic Beam; and U.S. Pat. No. 5,903,516, Acoustic Force Generator for Detection, Imaging and Information Transmission Using the Beat Signal of Multiple Intersecting Sonic Beams.

Experimental Embodiment

Referring to FIG. 3B, a possible AFM for deflecting a cantilever by focusing ultrasonic energy in the manner described above is illustrated at 110. The AFM 110 includes the probe 112 described above in conjunction in FIG. 3A (the probe 112 is configured to operate in a liquid contained within a cell 114), an ultrasonic actuator 116 for generating ARP, and a detector 118, 120. As discussed above, the probe 112 includes a cantilever 122 and a tip 124 mounted or otherwise provided on a free end of the cantilever 122. The cantilever of this example is made of silicon and has a spring constant of 0.148 N/m and a fundamental resonance around 4.6 kHz in water. The base of the cantilever 122 is fixed to an optically transparent probe support 126, which may fit into a commercial AFM scanhead 128. The fluid cell 114 is positioned beneath the scanhead 128 with the probe 112 located by it. (Relative positional terms such as "above", "beneath" "in front of" "behind," "horizontal," "vertical," etc. are used by way of reference for simplicity sake and are in no way limiting.) An ultrasonically transparent substrate 130, preferably made of a hard substrate like glass or silicon, is placed below the probe 112 and supports fluid cell 114. The ultrasonic actuator 116, which is preferably formed from a zinc oxide (ZnO) transducer, is mounted on the bottom of the substrate 130. Alternatively, the ultrasonic actuator can be placed on the side of the substrate directly facing the cantilever. The ultrasonic transducer can be shaped as a focusing device and electrically isolated from the fluid environment, removing the need for an ultrasonically transparent substrate. The ultrasonic actuator 116 is driven by a RF signal generator 132 to generate a beam 134 that deflects the cantilever 122 away from the substrate 130. The RF signal generator 132 has an optional modulation input that allows the amplitude of the RF signal to be varied with time. The modulation signal may be a square wave, a sinusoidal wave, a triangle wave or an arbitrary time-varying modulation. The RF signal generator also has an input or an internal adjustment that allows control over the baseline (unmodulated) power of the RF signal. The scanhead 128 may include an XY actuator and a Z actuator to permit the probe 112 to scan a sample in the fluid cell 114. Alternatively, the scanhead 128 could be stationary, and the substrate 130 could be driven to scan relatively to the scanhead 128. Detector 118, 120 detects cantilever deflection during scanning. The detector includes a laser 118 positioned above the cantilever 122 and a four-quadrant photodetector 120 configured to receive laser light reflected from the upper surface of the cantilever 122. As is conventional, signals from the photodetector 120 can be used as feedback to control operation of the RF signal generator 132 to maintain a desired characteristic of cantilever deflection, such as magnitude, and/or phase during scanning.

The substrate 130 of this embodiment contains a surface micromachined acoustic Fresnel lens 136 that serves as the focusing device in the illustrated embodiment. The illustrated lens structure is part of a 2-D array of micromachined acoustic lenses on the same glass plate. Lenses of this type were originally developed for acoustic microscopy and ink printing purposes and are well known. The Fresnel lens 136 of this embodiment is preferably designed such that, when the zinc oxide transducer 116 is excited with a sinusoidal drive signal at a frequency equal to about 179 MHz, the ultrasonic beam 134 will be focused to a diameter of approximately 5 μm to 10 μm at a focal distance of 360 μm . The 5 μm minimum diameter is even smaller than the 8 μm to 12 μm diameter of most laser beams. As a result, the lens 136 can be used to apply a pinpoint force to the free end of the cantilever 122 or any other point of interest along the length of the cantilever 122. Of course, the focal length of the Fresnel lens can be varied to accommodate any physical design constraints to place the ultrasonic actuator further or closer from the cantilever. Other types of acoustic lenses may also be used to shape the beam. A wide variety of acoustic lenses and beam shaping devices have been developed for medical ultrasound applications (phased arrays), scanning acoustic microscopy, acoustic printing and related techniques. For example, a simple hemispheric cutout in the surface of the substrate 130 will form an acoustic lens that will focus the outgoing beam. Further, materials with different speeds of sound may be patterned on top of the acoustic actuator to shape the profile of the outgoing beam. Further, the electrodes on the ultrasonic actuator may be patterned to form focused and/or steered beams using techniques such as Fresnel zone plates.

Characteristics of an Ultrasonically Driven Cantilever

Applying an RF drive voltage to the ZnO transducer 116 deflects the cantilever 122 in the manner illustrated in FIG. 4, which plots the RF signal input to the ultrasonic transducer and the resulting cantilever deflection via the curves 140 and 142, respectively. To obtain this data, the cantilever 122 of FIG. 3B was positioned on or near the focal plane of the ZnO transducer 116 using the Z-actuator of the scanhead 128. The X-Y actuator in the scanhead 128 was then driven to move the cantilever 122 in the X-Y (horizontal) plane to bring the tip 124 generally to the focal spot of the acoustic beam 134. The Z actuator was then driven to move the cantilever 122 vertically to obtain maximum cantilever deflection for a prevailing RF generator setting. The RF signal's amplitude was then modulated to 200 Hz at a frequency of about 179 MHz by a 50% duty cycle square wave with a 5 msec period. Then, after a transient dominated by its fundamental resonance in water, the cantilever 122 was driven upwardly away from the substrate 130, exhibiting a step response as represented by the curve 142 in FIG. 4. The particular response reflected by the curve 142 is consistent with the fact that the cantilever 122 of this exemplary embodiment has a fundamental resonance around 4.6 kHz in water. Corresponding data obtained from different cantilevers would produce curves of different magnitude and spacing but similar shape. The low frequency appearance in the RF drive signal curve 140 is due to the "aliasing" from the relatively low sampling rate of the digital oscilloscope (100 kS/s) used to obtain the data reflected in FIG. 4. The high signal-to-noise ratio obtained in the waveforms indicated by the curve 140 shows the potential of the method of surface characterization.

As should be apparent from above discussion, the force imposed by the ZnO transducer 116 is unidirectional. As a

result, the ultrasonic beam **134** cannot pull the cantilever **122** toward the transducer **116** but, instead, can only “push” the AFM cantilever **122** in the propagation direction. If the designer wanted to configure the AFM **110** to selectively pull the cantilever **122** toward the substrate **130** rather than push it away from the substrate, the cantilever **122** could be manufactured with a bias that maintains it in contact with the substrate **130** in the absence of a drive signal to the RF signal generator **132**. The RF signal generator **132** could then be driven to overcome the bias and push the cantilever **122** from the substrate **130**. The drive voltage could then be reduced to permit the cantilever **122** to move towards the substrate **130**, hence, in effect, pulling the cantilever **122** towards the substrate **130**. Alternatively, in some configurations, an ultrasonic actuator could be placed both above and below the cantilever to push from both sides.

The previous paragraph demonstrates a very important capability of the current invention. The ultrasonic cantilever actuator can independently control both the DC and AC forces applied to the cantilever over an extremely wide bandwidth. This opens up a large range of applications for this actuator. One example is an imaging method called Force Modulation. In this method, an AFM tip is brought into contact with a sample surface and then an AC modulation force is applied to the cantilever. The detector then measures the amount of AC deflection of the cantilever. On hard samples, the cantilever cannot indent into the surface, and no deflection is detected. On softer samples, the ultrasonic modulation force causes the tip to indent into the sample, resulting in a measurable AC deflection of the cantilever. Separate control over the DC force allows control over the tracking force that the AFM system uses to maintain contact between the tip and the surface.

The RF voltage can also be varied slowly to permit a quasistatic measurement to be performed. A quasistatic force imposition process is considered to be one in which, if the forces were to be removed at any stage during the process, system would be in equilibrium from that time on. Hence, the RF voltage can be changed slowly enough to maintain equilibrium while the voltage is being altered. This procedure is in contrast to a dynamic (AC) measurement in which the RF signal is modulated at a high frequency and the system requires time to stabilize when force adjustment terminates. The cutoff between quasistatic and dynamic measurements is usually considered to be a frequency value below the cantilever’s fundamental resonant frequency.

The force applied to the cantilever **122** is also proportional to the magnitude of the drive voltage applied to the RF signal generator **132**. This relationship is illustrated by FIG. **5**, the curve **146** of which traces measured cantilever deflection as a function of the input RF voltage to the ZnO transducer **116**. The cantilever used to obtain this data was a 205 μm thick, 2 μm thick, 29 μm wide, and had a spring constant of 1.14 N/m. As expected from Equation 2 above, the curve **146** indicates that the radiation force on the cantilever **122** varied generally linearly with input power to the transducer **116** over a wide operating range.

The generally linear relationship between drive voltage and cantilever deflection can be relied upon to calibrate the force applied to the cantilever for a particular drive voltage if the spring constant of a reference cantilever is known. For example, in the graph of FIG. **5**, a cantilever deflection of 31 nm was measured when 4.98 mW was input to the ZnO transducer. In an experiment using the data, the cantilever was assumed to have a spring constant of 1.14 N/m. That assumption would be verified if a force of 35 nN would have to be applied to the free end of the cantilever to deflect it. An

independent measurement of the insertion loss of the acoustic transducer and lens combination was performed to determine the incident acoustic power at the cantilever surface. This measurement yielded loss of 26.5 dB. Therefore, for the 4.98 mW electrical power on the ultrasonic transducer, the incident acoustic power on the tip was determined to be 11.1 μW . Using Equation 2, the force on the cantilever due to acoustic radiation pressure was predicted to be 13 nN, assuming an average reflection coefficient 0.85 for the 2 μm thick cantilever immersed in water. This number is in the same order of magnitude as the value predicted from the cantilever deflection. The calibrated force was then used to measure the spring constants of 405 and 105 μm long cantilevers on the same chip. The measured values were 0.132 N/m and 10.18 N/m, respectively. These values were 12% smaller and 19% larger, respectively, than the corresponding values quoted by the vendor.

The spring constant of a cantilever can be also determined without directly calibrating force by comparing the deflection of the cantilever at a particular drive voltage to the deflection of a reference cantilever of known spring constant at the same drive voltage. Specifically, the deflection of a reference cantilever of a known spring constant can be measured at one or more drive voltage(s), and the reference cantilever can then be replaced with one of an unknown spring constant. The same drive voltage(s) can be supplied to the RF signal generator **132**, and the cantilever’s deflection can be measured at the drive voltage(s). The cantilever’s spring constant can then be determined simply by determining the ratios of the two deflections at the same drive voltage.

Cantilever spring constants may also be determined dynamically, by ramping the ultrasonic force up and down. In this case, the spring constant of the cantilever is related to the slope of the cantilever deflection versus ultrasonic drive voltage. Since this is an AC measurement, this method has the advantage of being less sensitive to DC drift in the deflection of the cantilever or drift in the detection system due to changes in temperature.

The proportional relationship between cantilever deflection and drive voltage magnitude can also be used to obtain useful information concerning a sample. For instance, the RF signal generator **132** of FIG. **3B** can be controlled to generate force curves. Force curves are often used to provide an indication of the magnitude of force required to obtain an effect such as indenting a sample surface, breaking the binding molecules between a sample and a probe in contact with the sample, etc. In order to generate a force curve with prior instruments, it was necessary to actuate a Z actuator in the scanhead to drive the probe against the sample (or move it away from the sample). The same effect can be obtained using an ultrasonic actuator simply by modulating the drive signal to an RF signal generator or other power source until the desired effect is achieved. The known proportional relationship between the signal drive voltage and cantilever deflection can then be used to calculate the force. No actuation of a separate Z actuator is required.

AFM cantilevers can also respond dynamically to radiation pressure, and those dynamic responses can be measured. Specifically, time harmonic forces can be generated by applying a sinusoidal amplitude modulation on the RF input signal. By choosing the modulation factor to be less than one, an appropriate biasing force can be applied to actuate the cantilever at the modulation frequency and its second harmonic. The deflection of the cantilever can then be recorded e.g., by using a lock-in amplifier, which uses the modulation signal as its reference input and locks to the

modulation frequency. The normalized magnitude of the lock-in amplifier, was output as a function of the modulation frequency for a 405 μm calibration cantilever, a 205 μm calibration cantilever, and a V-shaped, diamond coated force modulation cantilever. That output is plotted in the curves **148**, **150**, and **152** respectively in FIG. 6. (The frequency sweep for this example was limited to 0.2–100 kHz range due to short time constant and the limitations of the lock-in amplifier.) The points **154** and **156** on curve **148** reveal that the fundamental and second modes of the long cantilever are about 4.6 kHz and 38 kHz, respectively. The points **158** and **160** on the curves **150** and **152** indicate that the short reference cantilevers and force modulation cantilever both resonate at about 25 kHz. The curves **148**, **150**, and **152** of FIG. 6 also show that the radiation pressure method of cantilever deflection can be used to actuate different cantilevers without any undesired effects of the cantilever support and the fluid cell. It has to be noted that, although only the flexural modes of the cantilevers are excited in this example, torsional modes of the cantilever can also be characterized by applying the radiation pressure at off-axis locations on the cantilever.

The relative ease with which both a resonance peak of an oscillation cantilever can be identified is confirmed with reference to FIG. 7A, which plots RMS deflection vs. drive frequency for a relatively short, thick cantilever (having a length of about 100 μm and a thickness of about 700 μm that is excited alternatively via an acoustic drive of the type illustrated in FIG. 1C and an ultrasonic drive of the type illustrated in FIG. 3B, respectively. The true response of the cantilever at rest, as determined via a known thermal tuning technique, is reflected by the curve **162**. Point **164** of curve **162** indicates that the cantilever actually experiences a fundamental resonance at about 15 kHz. Curve **166** indicates that it is virtually impossible to differentiate the cantilever resonance from other, parasitic resonances when the cantilever is driven by a conventional acoustic drive. In sharp contrast, the cantilever resonance can be clearly detected at point **170** in curve **168**. This effect is believed to be due at least in part to the fact that the ZnO transducer 1) acts along the length of the cantilever and, therefore has high gain, and 2) acts at least primarily on the cantilever rather than other components of the system and, accordingly, induces far fewer and smaller-amplitude parasitic resonances.

The measured response to ultrasonic excitation as reflected by the curve **168** in FIG. 7A is a true response at the excitation amplitude imparted by the ZnO transducer. The frequency shift between the peaks **164** and **170** of curves **162** and **168** results only from the known effects of finite excitation amplitude, not any measurement error. Hence, exciting the cantilever with an ultrasonic actuator permits the response of the cantilever to a given RF voltage to be anticipated with great confidence, even when the probe is operated in a liquid. Measurements can therefore be taken using cantilever/medium combinations that were heretofore not possible.

It has been discovered that an ultrasonic-actuator based system has a very wide bandwidth for exciting the cantilever. The inventors have performed experiments where the RF excitation signal is modulated at frequencies of more than 5 MHz, allowing unprecedented bandwidth for cantilever actuation as demonstrated by the curve **180** in FIG. 7B, which plots cantilever oscillation amplitude as a function of RF frequency at a 200 Hz modulation rate. In practice, it is possible to modulate the RF signal up to about $\frac{1}{10}$ of the RF signal frequency or even higher. For a 300 MHz RF frequency, a cantilever actuation bandwidth of even 30 MHz

is realizable. In fact, it is believed that the RF frequency from a low of 10 MHz or possibly even lower for air applications to 1 GHz or even higher for water applications with surface micromachined cantilevers for which attenuation may not be a problem. This cantilever actuation bandwidth is much greater than is provided by any other AFM cantilever actuators, particularly acoustic and magnetic drives of the type illustrated in FIG. 1C or ultrasonic force actuators of the type illustrated in FIG. 1D. An AFM having an ultrasonically driven cantilever can therefore be used to scan at rates that would have heretofore been considered outrageously fast. It also permits dynamic experiments to be conducted in which the probe repeatedly interacts with the sample surface at very high frequencies.

Tapping Mode Embodiment

An ultrasonic actuator of the type described above can be used to drive an AFM cantilever to oscillate at virtually any desired frequency significantly below the RF carrier frequency. An ultrasonic actuator therefore can be directly used as TappingMode actuator in an AFM. A TappingMode AFM **210** using an ultrasonic actuator is shown schematically in FIG. 8. As with the more theoretical embodiment of FIG. 3B, the instrument **210** includes a conventional probe **212**, an ultrasonic actuator **216**, and a detector **220**. The probe **212** is configured to operate in a fluid cell **214** containing a sample S. It includes a cantilever **222** having a base affixed to a support **226** and a free end bearing a tip **224**. Also as in the embodiment of FIG. 3B, an ultrasonically transmissive substrate **230** is placed below the cantilever **222**. A focusing device such as a Fresnel lens (not shown) may, if desired, be placed on or in the substrate **230**. The ultrasonic actuator **216** is mounted on the bottom of the substrate **230** and powered by an RF signal generator **232** under the control of an AFM controller **238**. The RF signal generator **232** produces an RF oscillation in the ultrasonic actuator **216** and then modulates the amplitude of that signal in response to a TappingMode drive signal as supplied by the controller **238**. The controller **238** may turn the drive signal on and off with a square wave such as the one illustrated in FIG. 4, or it could modulate the amplitude of the drive signal in proportion to that sine wave. The controller **238** also drives an XYZ actuator in the scanhead **228** in the conventional manner.

The traditional TappingMode piezoelectric drive may be taken out of the loop and replaced by the ultrasonic actuator **216**. In this case, the RF drive signal described in the preceding paragraph would always be used to drive the ultrasonic actuator **216**. In the preferred embodiment however, a piezoelectric drive **240** can be retained, and a suitable switch **242** can be provided to permit a drive signal to be selectively transmitted to either modulate the output of RF signal generator **232** and hence activate the ultrasonic actuator, or the drive signal can be sent to the piezoelectric drive **240** directly from the AFM controller **238**. For example, the resulting instrument could be operated in either air or liquid, with the piezoelectric drive **240** being used to effect operation in air and the ultrasonic actuator **216** being used to effect operation in liquid.

Overhead Ultrasonic Actuator Embodiment

A limitation of the instruments illustrated in FIGS. 3B and 8 is that the sample must be ultrasonically transmissive to permit unfettered transmission of the ultrasonic beam from the ultrasonic actuator, through the sample, and to the cantilever. A more versatile ultrasonic actuator assembly is schematically illustrated in FIG. 9. In this embodiment, at

least one ultrasonic actuator is mounted on a holder **330** positioned above the cantilever **322**. The holder **330** is also positioned between the cantilever **322**, on the one hand, and a laser **318** and photodetector (not shown), on the other hand. The holder **330** may be constructed of glass or any other material that is transparent to light and transmissive to ultrasonic energy. In the illustrated embodiment, two ZnO transducers **316a** and **316b** are mounted on the holder **330** on opposite sides of a vertical plane laterally bisecting the cantilever **322**, and the relevant portions of the holder are inclined to direct the corresponding ultrasonic beams **334a**, **334b** at the lateral centerline of the longitudinal centerline of the cantilever **322**. However, a single ultrasonic transducer would suffice in many applications. A Fresnel lens or other focusing device (not shown) could be formed in or mounted on the holder **330**, if desired. The holder, ultrasonic actuator (s), and associated components of this embodiment could also be used in place of the corresponding components of the embodiments of FIGS. 3B and 8, widening the range of applications of those embodiments.

Cantilever Characterization and Mode Shape Imaging

One of the unique features of an ultrasonic actuator is that it can produce a localized force at the desired location on the cantilever. It is therefore possible to very precisely control the vibration of a cantilever and to excite more flexural and torsional modes in the cantilever. These effects are illustrated in FIG. 10A, in which calculated and measured curves of the mode shape of a 450 μm cantilever in the AFM of FIG. 8 are shown at **350** and **352**, respectively. The data was taken at an excitation frequency of 1.5 kHz, which corresponds to the first resonance mode of this cantilever in water. As one would expect, the curves **350** and **352** reveal that cantilever deflection is maximized when the ultrasonic energy is directed at the free end of the cantilever and decreases when the focus point moves progressively toward the base. Stated another way, exciting the cantilever in its first resonance mode produces a single deflection peak.

A less intuitive characteristic of exciting a cantilever with a focused beam is that driving a cantilever to oscillate at its second and higher resonance modes produces a number of deflection peaks that increases with the order of the resonance mode. Hence, referring to the curves **354** and **356** of FIG. 10B, when the cantilever described in the preceding paragraph is excited to oscillate at its third resonance mode (58 kHz in the illustrated embodiment), three distinct deflection peaks are produced along the length of the cantilever. These peaks are denoted by points **358**, **360**, and **362**, in the curve **356**. Phase versus position along the cantilever is graphed by the curve **364** in FIG. 10B. Hence, cantilever deflection can be maximized or nearly maximized (hence maximizing cantilever response) by directing a beam at any of a number of different regions of the cantilever.

One important use of this capability is for surface elasticity characterization. Since the location of the nodes in these mode shapes are very sensitive to the surface properties at the tip-sample contact, the actuator can be positioned at a specific location relative to the cantilever free end and driven at a specific modulation frequency at which either the amplitude or phase of the mode shape changes very rapidly.

An instrument **410** configured for elasticity characterization is illustrated in FIG. 11. It includes all of the components of the AFM of FIG. 8, including probe **412**, a fluid cell **414**, an ultrasonic transducer **416**, a scanhead **428**, a substrate **430**, and a piezoelectric drive **440**. The probe **412**

includes a cantilever **422** that has a base mounted on a support **426** and that has a free end bearing a tip **424**. Electronic components of the instrument **410** include an RF signal generator **432**, an AFM controller **438**, and a switch **442** selectively coupling the controller **438** to the RF signal generator and the piezoelectric drive **440**. They additionally include a lock-in amplifier **444** and a low frequency signal generator **446**. The lock-in amplifier **444** receives a feedback signal from the low frequency signal generator **446** and transmits an elasticity image signal to the AFM controller **438**. Using conventional feedback to keep the force applied to the sample **S** constant, an image can be formed by monitoring cantilever deflection at the modulation frequency. The image, taking the form of an AC signal, has an amplitude and phase that both vary as a function of the sample stiffness. The procedure therefore can yield two simultaneous images: one for topography and one for surface elasticity. Elasticity characterizations using this technique can be performed much more rapidly than with prior known techniques due to the fact that the ultrasonic actuator has a dramatically higher bandwidth than prior systems that relied on a piezoelectric actuator to drive the entire probe up and down to obtain the required measurements. The prior systems typically had a maximum modulation frequency of only a few tens of kHz, whereas (as discussed above) an ultrasonic actuator based system has a cantilever actuation bandwidth of several MHz.

The same technique can be used to excite and measure the cantilever in torsional modes. For instance, FIGS. 12a and 12b show the amplitude and phase of the first two torsional modes of a V-shaped silicon nitride AFM cantilever.

No Focus Embodiment

A focusing device is not required at all if the ultrasonic actuator is configured to produce a beam that is collimated, minimally divergent, or otherwise configured to impinge on the cantilever with sufficient precision to negate the need for a focusing device. An ultrasonic actuator **516** configured to produce a collimated beam is illustrated in FIG. 13. The actuator **516**, which is mounted on a transparent holder **530**, is configured to produce a slowly diverging beam **534**, and the cantilever **522** is positioned with its body and free end in the beam **534** and the base and support **526** outside of the beam.

Beam divergence can be minimized using well-known ultrasonic transducer design techniques such as are disclosed, for example, in G. S. Kino "Acoustic Waves, devices, imaging and analog signal processing" Prentice-Hall, 1987, Englewood Cliffs, N.J. In general, the requirements for minimum beam divergence can be derived from diffraction calculations. One way of optimizing the energy collimation would be to place the probe at a distance where the near-field to far-field transition happens. For a circular actuator with diameter D , this distance is given by $L=D^2/\lambda$, where D is the diameter of the actuator and λ is the wavelength of the ultrasonic waves in the medium (λ =speed of sound/frequency). For a rectangular actuator which is very long in one dimension (as shown in FIG. 9) and has of width W , the optimum distance will be $L=W^2/4\lambda$. At this distance, the ultrasonic energy will be dominantly on the axis. For smaller distances, in the near-field, the ultrasonic energy will rapidly fluctuate, and the exact position of the probe will be more critical.

For distances larger than $L=D^2/\lambda(W^2/4\lambda)$, the beam will diverge with a divergence angle of $\alpha=\sin(1.22\lambda/D)$, for a circular actuator and $\alpha=\sin(\lambda/W)$ for a rectangular trans-

ducer. Therefore, one way (but by no means the only way) of limiting beam divergence is by maximizing the frequency f of the beam **534** and minimizing the diameter D at the base of the beam as determined by the diameter of the ultrasonic actuator **516**. Because f is inversely proportional to the wavelength (λ) of the ultrasonic signal, the divergence (α) can be kept very small by minimizing both D and λ within practical limits. (As discussed above, this may not be the case for near field applications) A divergence of less than 10 degrees is preferred, but an effective actuator could still be designed with a divergence of 30 degrees or more. This effect can be achieved by placing a small ZnO actuator having a mean width on the order of 50–500 microns directly above the cantilever at a spacing of L up to several mm without any focusing. The actuator can be made in a variety of shapes—circles, ovals, rectangles or other arbitrary shape to provide the desired ultrasonic beam profile.

Using either focusing or collimation, the beam may be made intentionally smaller than the cantilever so that all of the energy strikes the cantilever. This embodiment is preferred for applications where the ultrasonic actuator is used to apply a very well known force to the cantilever, for example to measure the spring constant of a cantilever or to apply a known force from the cantilever to the sample. In an alternate embodiment, the beam may be intentionally sized larger than the cantilever to account for tolerances in the alignment of the cantilever and the ultrasonic actuator. If for example, the cantilever is 50 microns wide and can be reproducibly aligned within ± 100 microns, an ultrasonic actuator with a beam width of 250–300 μm in the region of the cantilever could guarantee that a portion of the ultrasonic beam would always strike the cantilever. In the preferred embodiment, actuator sizes range from minimum widths of about 50 μm up to about 3 mm.

AFM Array with Integrated ARP Actuation

The ARP method can be easily used to actuate AFM cantilevers in an array. FIG. **14A** shows the schematic of one of these cantilever arrays. A plurality of ultrasonic actuators **616a**, **616b**, etc., are mounted on a common transparent holder **630**. Each actuator is configured to direct a beam (beam **634d** for example) of ultrasonic energy at a corresponding cantilever **622a**, **622b**, etc. of an array of cantilevers mounted on a common support **624**. As with the previous embodiments, the actuators can be accompanied by focusing lenses, collimation or other beam confinement, as necessary.

In an alternate embodiment shown in FIG. **14B** a single rectangular aperture can be provided in a holder **730** and configured to transmit a wide ultrasonic beam **734** from a single ultrasonic actuator **716** that strikes all the cantilevers **722a**, **722b**, etc. of the array at once. The divergence of the beam **734** can be reduced as necessary using a lens or using the techniques described above in conjunction with the “No Focus” embodiment of FIG. **13**. Cantilever arrays for AC applications are usually designed to have different resonant frequencies so that the oscillation of one cantilever does not excite a sympathetic vibration of adjacent cantilevers. With this type of array, the modulation frequency can be tuned to oscillate one cantilever at a time so that the dynamic properties of the individual cantilevers in the array can be measured. The rectangular actuator arrangement **716** of FIG. **14B** requires higher power to energize the larger actuator, but it has the advantage of not requiring a switching circuit to energize the actuators individually as in FIG. **14A**.

Although the best mode contemplated by the inventors of carrying out the present invention is disclosed above, prac-

tice of the present invention is not limited thereto. It will be manifest that various additions, modifications and rearrangements of the features of the present invention may be made without deviating from the spirit and scope of the underlying inventive concept. For instance, and as mentioned briefly above, although the examples described above focus on imaging in liquids such as water, the invention is also applicable to imaging in less dense fluids and even to gases such as air. The scope of still other changes to the described embodiments that fall within the present invention but that are not specifically discussed above will become apparent from the appended claims.

We claim:

1. A metrology instrument comprising:

(A) a probe including a cantilever; and

(B) an ultrasonic actuator configured to direct a beam of ultrasonic energy at the cantilever that imposes a force on the cantilever that drives the cantilever to oscillate.

2. The instrument as recited in claim 1, wherein the ultrasonic actuator is configured to drive the cantilever to oscillate substantially at the resonant frequency of the cantilever or a harmonic thereof.

3. The instrument as recited in claim 2, wherein at least a free end portion of the cantilever is configured to be submerged in a liquid.

4. The instrument as recited in claim 1, wherein the instrument is configured to transmit an RF oscillation signal to the ultrasonic actuator, and wherein the amplitude of the RF oscillation signal is adjustable to provide an adjustable force to the cantilever.

5. The instrument as recited in claim 4, wherein the ultrasonic transducer has a modulation frequency bandwidth of at least 1 MHz.

6. The instrument as recited in claim 5, wherein the ultrasonic transducer has a modulation frequency bandwidth of at least 5 MHz.

7. The instrument as recited in claim 4, wherein the ultrasonic transducer has a modulation frequency bandwidth of at least 50 kHz.

8. The instrument as recited in claim 4, further comprising a second actuator configured to drive the cantilever to oscillate, and wherein the instrument is switchable between a first state in which an oscillating drive voltage is supplied to modulate the RF excitation to the ultrasonic actuator and a second state in which an oscillating drive voltage is supplied to the second actuator.

9. The instrument as recited in claim 1, further comprising a detector that is configured to detect cantilever deflection.

10. The instrument as recited in claim 9, wherein the cantilever and the detector are positioned on a common side of the cantilever disposed opposite a sample holder, and wherein the ultrasonic actuator comprises a first ultrasonic actuator, and further comprising a second ultrasonic actuator configured to direct a beam of ultrasonic energy at the cantilever, wherein the first and second ultrasonic actuators flank a longitudinal centerline of the cantilever.

11. The instrument as recited in claim 9, wherein the detector comprises a photodetector configured to detect a light beam reflected from the cantilever.

12. The instrument as recited in claim 1, wherein the cantilever comprises one of an array of cantilevers, and wherein multiple ultrasonic actuators are provided for directing beams of ultrasonic energy at multiple cantilevers of the array.

13. The instrument as recited in claim 1, wherein the cantilever comprises one of an array of cantilevers, and wherein a single ultrasonic actuator is provided for directing a beam of ultrasonic energy at one or more cantilevers of the array.

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14. A metrology instrument comprising:

(A) a cantilever;

(B) an ultrasonic actuator configured to direct a beam of ultrasonic energy at the cantilever so as to drive the cantilever to oscillate;

(C) an RF excitation signal; and

(D) a modulation signal generator configured to generate a modulation signal that modulates the amplitude of the RF excitation signal to produce a modulated RF signal that is transmitted to the ultrasonic actuator.

15. The instrument as recited in claim 13, wherein the modulated RF signal has a bandwidth of at least 50 kHz.

16. The instrument as recited in claim 14, wherein the modulated RF signal has a bandwidth of at least 1 MHz.

17. A scanning probe microscope (SPM) comprising:

(A) a probe including a cantilever;

(B) an ultrasonic actuator coupled to the probe so as to drive the cantilever in a selected mode of SPM operation, wherein, when driving the cantilever, the frequency response of the cantilever has amplitude peaks at substantially only the fundamental resonance of the cantilever and harmonics of the fundamental resonance.

18. A method comprising:

(A) generating a beam of ultrasonic energy using an ultrasonic actuator;

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(B) transmitting the beam onto a cantilever of a probe of a metrology instrument to impose a force on the cantilever so as to drive the cantilever to oscillate.

19. The method as recited in claim 18, wherein the cantilever is driven to oscillate at the resonant frequency of the cantilever or a harmonic thereof.

20. The method as recited in claim 18, further comprising adjusting a drive frequency supplied to the ultrasonic actuator through a range at least 50 kHz.

21. The method as recited in claim 19, further comprising focusing the beam onto a region of the cantilever so as to produce peak deflection when the cantilever is oscillated in a second or higher order resonance mode.

22. The method as recited in claim 17, wherein at least a free end of the cantilever is immersed in a liquid during the transmitting step, and further comprising reconfiguring the cantilever so that it is not immersed in a liquid and driving the cantilever to oscillate using a second actuator.

23. The method as recited in claim 17, further comprising detecting cantilever deflection.

24. The method as recited in claim 22, further comprising taking elasticity measurements using data collected as a result of the detecting step.

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